#### Quantum speedup of Monte Carlo methods

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Pic: Wikipedia

#### These methods are used throughout science and engineering:



... and were an application of the first electronic computers:



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One very simple example of a Monte Carlo method: approximate  $\pi$  by throwing darts at a dartboard (choosing random points within a square).



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## A randomised algorithm for approximating $\pi$

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#### This talk

Today I will discuss a quantum algorithm to speed up Monte Carlo methods in a quite general setting.

And also some applications of the algorithm:

- Partition function problems in statistical physics
- **2** Approximate counting problems in combinatorics
- Approximating the distance between probability distributions

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- The input is fixed, and the expectation is taken over the internal randomness of *A*.
- The output value v(A) is a real-valued random variable.

We assume that we know an upper bound on the variance of this random variable:

 $\mathsf{Var}(v(\mathcal{A}))\leqslant\sigma^2.$ 

The following natural algorithm solves this problem for any A:

- Produce *k* samples  $v_1, \ldots, v_k$ , each corresponding to the output of an independent execution of A.
- Output the average  $\tilde{\mu} = \frac{1}{k} \sum_{i=1}^{k} v_i$  of the samples as an approximation of  $\mu$ .

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• To estimate  $\pi$  up to 4 decimal places with success probability 0.5, we would need > 10<sup>9</sup> trials!

With a quantum computer, we can do better:

Theorem [AM '15]

There is a quantum algorithm which estimates  $\mu$  up to additive error  $\varepsilon$  with 99% success probability and

 $\widetilde{O}(\sigma/\varepsilon)$ 

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The underlying algorithm A can now be quantum itself.

This problem connects to several previous works, e.g.:

Approximating the mean of an arbitrary bounded function (with range [0, 1]), with respect to the uniform distribution. Quantum complexity: O(1/ε) [Heinrich '01], [Brassard et al. '11].

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- Estimating the expected value tr(*A*ρ) of certain observables *A* which are bounded [Wocjan et al. '09], or whose tails decay quickly [Knill et al. '07].

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Here we generalise these by approximating the mean output value of arbitrary quantum algorithms, given only a bound on the variance.

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- Then we replace the end of A with the map

$$|x\rangle|0\rangle\mapsto|x\rangle(\sqrt{1-\varphi(x)}|0\rangle+\sqrt{\varphi(x)}|1\rangle).$$

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- We can use amplitude estimation to approximate μ up to additive error ε, using A (and A<sup>-1</sup>) O(1/ε) times.

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In this case (using ideas of [Heinrich '01]):

- Divide up the output values of  $\mathcal{A}$  into blocks, such that in the *t*'th block  $2^{t-1} \leq v(\mathcal{A}) \leq 2^t$ .
- Use Õ(1/ε) iterations of the previous algorithm to estimate the average output values in each of the first O(log 1/ε) blocks, each divided by 2<sup>t</sup>.
- Sum up the results (after rescaling them again).



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- Sum up the results (after rescaling them again).

The constraint that  $\mathbb{E}[v(\mathcal{A})^2] = O(1)$  implies that the overall error is at most  $\epsilon$ .

The final step is to change the dependence on  $\mathbb{E}[v(\mathcal{A})^2]$  to a dependence on

$$\mathsf{Var}(v(\mathcal{A})) = \mathbb{E}[(v(\mathcal{A}) - \mu)^2] \leqslant \sigma^2.$$

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- Estimate the positive and negative parts separately.

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A similar idea works to estimate  $\mu$  up to relative error  $\epsilon$ : if  $\sigma^2/\mu^2 \leq B$ , we can estimate  $\mu$  up to additive error  $\epsilon \mathbb{E}[v(\mathcal{A})]$  with  $\widetilde{O}(B/\epsilon)$  uses of  $\mathcal{A}$ .

Consider a (classical) physical system which has state space  $\Omega$ , and a Hamiltonian  $H : \Omega \to \mathbb{R}$  specifying the energy of each configuration  $x \in \Omega$ . Assume that H takes integer values in the set  $\{0, \ldots, n\}$ .

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Encapsulates some interesting problems:

- Physics: The Ising and Potts models
- Computer science: counting *k*-colourings of graphs, counting matchings (monomer-dimer coverings), ...

- *A* := |Ω| can be exponentially large and *Z*(β) can be hard to compute; e.g. #P-hard. So we resort to randomised methods for approximating *Z*(β).
- We want to approximate Z(β) up to relative error ε, i.e. output Z̃ such that

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- One approach: multi-stage Markov chain Monte Carlo (e.g. [Valleau and Card '72, Stefankovič et al. '09]).

#### Multiple-stage Markov chain Monte Carlo

The basic framework of these methods:

- Let a cooling schedule be a sequence of inverse temperatures 0 = β<sub>0</sub> < β<sub>1</sub> < · · · < β<sub>ℓ</sub> = β.
- Express  $Z(\beta_{\ell})$  as the telescoping product

$$Z(\beta_{\ell}) = Z(\beta_0) \frac{Z(\beta_1)}{Z(\beta_0)} \frac{Z(\beta_2)}{Z(\beta_1)} \dots \frac{Z(\beta_{\ell})}{Z(\beta_{\ell-1})}.$$

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• Define random variables  $Y_i$  where  $\mathbb{E}[Y_i] = Z(\beta_{i+1})/Z(\beta_i)$ , with respect to the distribution  $\pi_i$  defined by

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• Estimate  $\mathbb{E}[Y_i]$  by sampling from this distribution.

# Sampling and estimating

This idea will be efficient if we can satisfy two constraints:

- The (relative) variance of each random variable  $Y_i$  is low:  $\mathbb{E}[Y_i^2]/\mathbb{E}[Y_i]^2 = O(1)$  for all *i*.
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#### Theorem [Stefankovič et al. '09]

For any partition function problem, there exists a cooling schedule satisfying the first constraint with  $\ell = \widetilde{O}(\sqrt{\log A})$ .

Such a cooling schedule is known as a Chebyshev cooling schedule.

# **Rapid mixing**

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- In the quantum setting, we can apply our algorithm to accelerate the approximation of E[Y<sub>i</sub>] (scaling goes from O(1/ε<sup>2</sup>) to Õ(1/ε))...
- ... and we can also replace the classical Markov chains with quantum walks to improve the dependence on τ.

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Note 2: The  $O((\log A)\tau)$  part of the bound is the complexity of computing the Chebyshev cooling schedule itself.

#### **Example: The ferromagnetic Ising model**

We are given as input a graph G = (V, E) with *n* vertices. We consider the Ising Hamiltonian

$$H(z) = -\sum_{(u,v)\in E} z_u z_v.$$

for  $z \in \{\pm 1\}^n$ . We want to approximate

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- Assume that we have a classical Markov chain which rapidly samples from the Gibbs distribution  $(\tau = \tilde{O}(n))$ .
- This holds for low enough β (depending on the graph *G*). Then we have the following speedup:
  - Best classical runtime known [Stefankovič et al. '09]:  $\tilde{O}(n^2/\epsilon^2)$
  - Quantum runtime:  $\widetilde{O}(n^{3/2}/\epsilon + n^2)$

# Applications

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Counting matchings (monomer-dimer coverings) of a graph with *n* vertices and *m* edges:

- Best classical runtime known:  $\tilde{O}(n^2m/\epsilon^2)$
- Quantum runtime:  $\widetilde{O}(n^{3/2}m^{1/2}/\epsilon + n^2m)$

- Imagine we can sample from probability distributions *p* and *q* on *n* elements.
- We would like to estimate the total variation distance

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- Using quantum mean estimation we improve this to  $\tilde{O}(\sqrt{n}/\epsilon^{3/2})$ .

• We can write  $||p - q|| = \mathbb{E}_x[R(x)]$ , where

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- It's sufficient to use  $O(\sqrt{n/\epsilon})$  iterations of amplitude estimation to approximate  $\mathbb{E}_x[R(x)]$  up to additive error  $\epsilon$ .
- Wrapping this within O(1/ε) iterations of the mean-estimation algorithm, we obtain an overall algorithm running in time Õ(√n/ε<sup>3/2</sup>).

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#### Thanks!