# The two-point correlation function of the fractional parts of $\sqrt{n}$ is Poisson<sup>\*</sup>

Daniel El-Baz<sup>†</sup>

Jens Marklof<sup>†</sup>

Ilva Vinogradov<sup>†</sup>

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#### Abstract

Elkies and McMullen [Duke Math. J. 123 (2004) 95–139] have shown that the gaps between the fractional parts of  $\sqrt{n}$  for  $n = 1, \ldots, N$ , have a limit distribution as N tends to infinity. The limit distribution is non-standard and differs distinctly from the exponential distribution expected for independent, uniformly distributed random variables on the unit interval. We complement this result by proving that the two-point correlation function of the above sequence converges to a limit, which in fact coincides with the answer for independent random variables. We also establish the convergence of moments for the probability of finding r points in a randomly shifted interval of size 1/N. The key ingredient in the proofs is a non-divergence estimate for translates of certain non-linear horocycles.

#### **1** INTRODUCTION

It is well known that, for every fixed  $0 < \alpha < 1$ , the fractional parts of  $n^{\alpha}$  (n = 1, ..., N) are, in the limit of large N, uniformly distributed mod 1. Numerical experiments suggest that the gaps in this sequence converge to an exponential distribution as  $N \to \infty$ , which is the distribution of waiting times in a Poisson process, cf. Fig. 1. The only known exception is the case  $\alpha = 1/2$ . Here Elkies and McMullen [2] proved that the limit gap distribution exists and is given by a piecewise analytic function with a power-law tail (Fig. 2). In the present study we show that a closely related local statistics, the two-point correlation function, has a limit which in fact is consistent with the Poisson process, see Fig. 3. The proof of this claim follows closely our discussion in [1], which produced an analogous result for the two-point statistics of directions in an affine lattice. (We note that Sinai [14] has recently proposed an alternative approach to the statistics of  $\sqrt{n} \mod 1$ , but will not exploit this here.)

Other number-theoretic sequences, whose two-point correlations are Poisson, include the values of positive definite quadratic forms subject to certain diophantine conditions

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<sup>&</sup>lt;sup>†</sup>School of Mathematics, University of Bristol, Bristol BS8 1TW, U.K.

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**Figure 1:** Gap distribution of the fractional parts of  $n^{1/3}$  with  $n \leq 2 \times 10^5$ .



**Figure 2:** Gap distribution of the fractional parts of  $\sqrt{n}$  with  $n \leq 2 \times 10^5$ .



**Figure 3:** Two-point correlations of the fractional parts of  $\sqrt{n}$  with  $n \leq 2000, n \notin \Box$ .

[12, 3], forms in more variables [16, 15, 17], inhomogeneous forms in two [8, 5] and more variables [7], and the fractional parts of  $n^2\alpha$  (and higher polynomials) for almost all  $\alpha$  [11, 6, 4] (specific examples, e.g.  $\alpha = \sqrt{2}$  are still open).

To describe our results, let us first note that  $\sqrt{n} = 0 \mod 1$  if and only if n is a perfect square. We will remove this trivial subsequence and consider the set

$$\mathcal{P}_T = \{\sqrt{n} \bmod 1 \colon 1 \leqslant n \leqslant T, \ n \notin \Box\} \subset \mathbb{T} := \mathbb{R}/\mathbb{Z}$$
(1.1)

where  $\Box \subset \mathbb{N}$  denotes the set of perfect squares. The cardinality of  $\mathcal{P}_T$  is  $N(T) = T - \lfloor \sqrt{T} \rfloor$ . We label the elements of  $\mathcal{P}_T$  by  $\alpha_1, \ldots, \alpha_{N(T)}$ . The pair correlation density of the  $\alpha_j$  is defined by

$$R_{N(T)}^{2}(f) = \frac{1}{N(T)} \sum_{m \in \mathbb{Z}} \sum_{\substack{i,j=1\\i \neq j}}^{N(T)} f(N(T)(\alpha_{i} - \alpha_{j} + m)), \qquad (1.2)$$

where  $f \in C_0(\mathbb{R})$  (continuous with compact support). Note that  $R^2_{N(T)}$  is not a probability density. Our first result establishes that  $R^2_{N(T)}$  converges weakly to the two-point density of a Poisson process:

**Theorem 1.** For any  $f \in C_0(\mathbb{R})$ ,

$$\lim_{T \to \infty} R_{N(T)}^2(f) = \int_{\mathbb{R}} f(s) \, ds.$$
(1.3)

Both the convergence of the gap distribution and of the two-point correlations follow from a more general statistics, namely, the probability of finding r elements  $\alpha_j$  in randomly placed intervals of size proportional to 1/N(T). Given a bounded interval  $I \subset \mathbb{R}$ , define the subinterval  $J = J_{N(T)}(I, \alpha) = N(T)^{-1}I + \alpha + \mathbb{Z} \subset \mathbb{T}$  of length  $N(T)^{-1}|I|$ , and let

$$\mathcal{N}_T(I,\alpha) = \#\mathcal{P}_T \cap J_{N(T)}(I,\alpha). \tag{1.4}$$

It is proved in [2] that, for  $\alpha$  uniformly distributed in  $\mathbb{T}$  with respect to the Lebesgue measure  $\lambda$ , the random variable  $\mathcal{N}_T(I, \alpha)$  has a limit distribution E(k, I). That is to say, for every  $k \in \mathbb{Z}_{\geq 0}$ ,

$$\lim_{T \to \infty} \lambda(\{\alpha \in \mathbb{T} : \mathcal{N}_T(I, \alpha) = k\}) = E(k, I).$$
(1.5)

As Elkies and McMullen point out, these results hold in fact for several test intervals  $I_1, \ldots, I_m$ :

**Theorem 2** (Elkies and McMullen [2]). Let  $I = I_1 \times \cdots \times I_m \subset \mathbb{R}^m$  be a bounded box. Then there is a probability distribution  $E(\cdot, I)$  on  $\mathbb{Z}_{\geq 0}^m$  such that, for any  $\underline{k} = (k_1, \ldots, k_m) \in \mathbb{Z}_{\geq 0}^m$ 

$$\lim_{T \to \infty} \lambda(\{\alpha \in \mathbb{T} : \mathcal{N}_T(I_1, \alpha) = k_1, \dots, \mathcal{N}_T(I_m, \alpha) = k_m\}) = E(\underline{k}, I).$$
(1.6)

In other words, Theorem 2 states that the point process

$$\{N(T)(\alpha_j - \alpha + \mathbb{Z})\}_{j \leq N(T)}$$

on the torus  $\mathbb{R}/(N(T)\mathbb{Z})$  converges, as  $T \to \infty$ , to a random point process on  $\mathbb{R}$  which is determined by the probabilities  $E(\underline{k}, I)$ . As pointed out in [10], this process is the same as for the directions of affine lattice points with irrational shift. It is described in terms of a random variable in the space of affine lattices and is in particular not a Poisson process. The second moments and two-point correlation function, however, coincide with those of a Poisson process with intensity 1. This is a consequence of the Siegel integral formula, see [1] for details (in the notation of [1],  $E(\underline{k}, I) = E_0(\underline{k}, I) = E_{0,\boldsymbol{\xi}}(\underline{k}, I)$  with  $\boldsymbol{\xi} \notin \mathbb{Q}^2$ ). Specifically, we have

$$\sum_{k=0}^{\infty} k^2 E(k, I_1) = |I_1| + |I_1|^2$$
(1.7)

and

$$\sum_{\underline{k}\in\mathbb{Z}_{\geq 0}^{2}}k_{1}k_{2}E(\underline{k},I_{1}\times I_{2}) = |I_{1}\cap I_{2}| + |I_{1}||I_{2}|.$$
(1.8)

The third and higher moments diverge [1].

It is important to note that Elkies and McMullen considered the full sequence  $\{\sqrt{n} \mod 1 : 1 \leq n \leq T\}$ . Removing the perfect squares  $n \in \Box$  does not have any effect on the limit distribution in Theorem 2, since the set of  $\alpha$  for which  $\mathcal{N}_T(I, \alpha)$  is different has vanishing Lebesgue measure as  $T \to \infty$ . In the case of the second and higher moments, however, the removal of perfect squares will make a difference and in particular avoid trivial divergence.

For  $I = I_1 \times \cdots \times I_m \subset \mathbb{R}^m$  and  $\underline{s} = (s_1, \ldots, s_m) \in \mathbb{C}^m$  let

$$\mathbb{M}(T,\underline{s}) := \int_{\mathbb{T}} (\mathcal{N}_T(I_1,\alpha) + 1)^{s_1} \cdots (\mathcal{N}_T(I_m,\alpha) + 1)^{s_m} d\alpha.$$
(1.9)

The main objective of the present paper is to establish the convergence of these mixed moments to the corresponding moments of the limit process, where they exist. The case of the second mixed moment implies, by a standard argument, the convergence of the two-point correlation function stated in Theorem 1, cf. Appendix 1 of [1].

We denote the positive real part of  $z \in \mathbb{C}$  by  $\operatorname{Re}_+(z) := \max\{\operatorname{Re}(z), 0\}$ , and set  $\sigma := \sum_{j=1}^m \operatorname{Re}_+(s_j)$ .

**Theorem 3.** Let  $I = I_1 \times \cdots \times I_m \subset \mathbb{R}^m$  be a bounded box, and  $\lambda$  a Borel probability measure on  $\mathbb{T}$  with continuous density. Choose  $\underline{s} = (s_1, \ldots, s_m) \in \mathbb{C}^m$ , such that  $\sigma < 3$ . Then,

$$\lim_{T \to \infty} \mathbb{M}(T, \underline{s}) = \sum_{\underline{k} \in \mathbb{Z}_{\geq 0}^m} (k_1 + 1)^{s_1} \cdots (k_m + 1)^{s_m} E(\underline{k}, I).$$
(1.10)

Our techniques permit to generalize the above results in two ways:

**Remark 1.** Instead of  $\mathcal{P}_T$  we may consider

$$\mathcal{P}_{T,c} = \{\sqrt{n} \bmod 1 \colon c^2 T < n \leqslant T, \ n \notin \Box\}$$
(1.11)

for any  $0 \leq c < 1$ . This setting is already discussed in [2, Section 3.5], and the upper bounds obtained in the present paper are sufficient to establish Theorem 3 in this case. Note that the limit process is different for each c; it coincides with the limit process  $E_c(\underline{k}, I)$ studied in [1]. As we point out in [1], the second moments of  $E_c(\underline{k}, I)$  are Poisson, and hence Theorem 1 holds independently of the choice of c.

**Remark 2.** Although Elkies and McMullen assume that  $\lambda$  is the Lebesgue measure, the equidistribution result that is used to prove Theorem 2, in fact, holds for any Borel probability measure  $\lambda$  on  $\mathbb{T}$  which is absolutely continuous with respect to Lebesgue measure. This follows from Ratner's theorem by arguments similar to those used by Shah [13]. It is important to note that the limit process  $E_c(\underline{k}, I)$  will be independent of the choice of  $\lambda$ . Theorem 3 then follows from the general version of Theorem 2 for measures  $\lambda$  with continuous density (since in this case, for all upper bounds, it is sufficient to restrict the attention to Lebesgue measure). As discussed in [1], the generalization of the above results to  $\lambda$  with continuous density yields the convergence of a more general two-point correlation function,

$$R_{N(T)}^{2}(f) = \frac{1}{N(T)} \sum_{\substack{m \in \mathbb{Z} \\ i \neq j}} \sum_{\substack{i,j=1 \\ i \neq j}}^{N(T)} f(\alpha_{i}, \alpha_{j}, N(T)(\alpha_{i} - \alpha_{j} + m)),$$
(1.12)

to the Poisson limit. That is, for all  $f \in C_0(\mathbb{T}^2 \times \mathbb{R})$ ,

$$\lim_{T \to \infty} R_{N(T)}^2(f) = \int_{\mathbb{T} \times \mathbb{R}} f(\alpha, \alpha, s) \, d\alpha \, ds.$$
(1.13)

#### 2 Strategy of proof

The proof of Theorem 3 follows our strategy in [1]. We define the restricted moments

$$\mathbb{M}^{(K)}(T,\underline{s}) := \int_{\max_j \mathcal{N}_T(I_j,\alpha) \leqslant K} (\mathcal{N}_T(I_1,\alpha) + 1)^{s_1} \cdots (\mathcal{N}_T(I_m,\alpha) + 1)^{s_m} d\alpha.$$
(2.1)

Theorem 2 implies that, for any fixed  $K \ge 0$ ,

$$\lim_{T \to \infty} \mathbb{M}^{(K)}(T, \underline{s}) = \sum_{\substack{\underline{k} \in \mathbb{Z}_{\geq 0}^m \\ |\underline{k}| \leqslant K}} (k_1 + 1)^{s_1} \cdots (k_m + 1)^{s_m} E(\underline{k}, I),$$
(2.2)

where  $|\underline{k}|$  denotes the maximum norm of  $\underline{k}$ . To prove Theorem 3, what remains is to show that

$$\lim_{K \to \infty} \limsup_{T \to \infty} \left| \mathbb{M}(T, \underline{s}) - \mathbb{M}^{(K)}(T, \underline{s}) \right| = 0.$$
(2.3)

To establish the latter, we use the inequality

$$\left|\mathbb{M}(T,\underline{s}) - \mathbb{M}^{(K)}(T,\underline{s})\right| \leq \int_{\mathcal{N}_{T}(\overline{I},\alpha) \geq K} (\mathcal{N}_{T}(\overline{I},\alpha) + 1)^{\sigma} d\alpha$$
(2.4)

where  $\overline{I} = \bigcup_j I_j$  and  $\sigma = \sum_j \operatorname{Re}_+(s_j)$ . As in the work of Elkies and McMullen, the integral on the right hand side can be interpreted as an integral over a translate of a non-linear horocycle in the space of affine lattices. The main difference is that now the test function is unbounded, and we require an estimate that guarantees there is no escape of mass as long as  $\sigma < 3$ . This means that

$$\lim_{K \to \infty} \limsup_{T \to \infty} \int_{\mathcal{N}_T(\overline{I}, \alpha) \ge K} (\mathcal{N}_T(\overline{I}, \alpha) + 1)^{\sigma} d\alpha = 0$$
(2.5)

implies Theorem 3. The remainder of this paper is devoted to the proof of (2.5).

### 3 ESCAPE OF MASS IN THE SPACE OF LATTICES

Let  $G = \mathrm{SL}(2,\mathbb{R})$  and  $\Gamma = \mathrm{SL}(2,\mathbb{Z})$ . Define the semi-direct product  $G' = G \ltimes \mathbb{R}^2$  by

$$(M,\boldsymbol{\xi})(M',\boldsymbol{\xi}') = (MM',\boldsymbol{\xi}M' + \boldsymbol{\xi}'), \qquad (3.1)$$

and let  $\Gamma' = \Gamma \ltimes \mathbb{Z}^2$  denote the integer points of this group. In the following, we will embed G in G' via the homomorphism  $M \mapsto (M, \mathbf{0})$  and identify G with the corresponding subgroup in G'. We will refer to the homogeneous space  $\Gamma \setminus G$  as the space of lattices and  $\Gamma' \setminus G'$  as the space of affine lattices. A natural action of G' on  $\mathbb{R}^2$  is defined by  $\mathbf{x} \mapsto \mathbf{x}(M, \boldsymbol{\xi}) := \mathbf{x}M + \boldsymbol{\xi}$ .

Given an interval  $I \subset \mathbb{R}$ , define the triangle

$$\mathfrak{C}(I) = \{ (x, y) \in \mathbb{R}^2 : 0 < x < 2, \ y \in 2|x|I \}.$$
(3.2)

and set, for  $g \in G'$  and any bounded subset  $\mathfrak{C} \subset \mathbb{R}^2$ ,

$$\mathcal{N}(g,\mathfrak{C}) = \#(\mathfrak{C} \cap \mathbb{Z}^2 g). \tag{3.3}$$

By construction,  $\mathcal{N}(\cdot, \mathfrak{C})$  is a function on the space of affine lattices,  $\Gamma' \setminus G'$ .

Let

$$\Phi^t = \begin{pmatrix} e^{-t/2} & 0\\ 0 & e^{t/2} \end{pmatrix}, \qquad \tilde{n}(u) = \begin{pmatrix} \begin{pmatrix} 1 & u\\ 0 & 1 \end{pmatrix}, \begin{pmatrix} \frac{u}{2}, \frac{u^2}{4} \end{pmatrix} \end{pmatrix}.$$
(3.4)

Note that  $\{\Phi^t\}_{t\in\mathbb{R}}$  and  $\{\tilde{n}(u)\}_{u\in\mathbb{R}}$  are one-parameter subgroups of G'. Note that  $\Gamma'\tilde{n}(u+2) = \Gamma'\tilde{n}(u)$  and hence  $\Gamma'\{\tilde{n}(u)\}_{u\in[-1,1)}\Phi^t$  is a closed orbit in  $\Gamma'\backslash G'$  for every  $t\in\mathbb{R}$ .

**Lemma 4.** Given an interval  $I \subset \mathbb{R}$ , there is  $T_0 > 0$  such that for all  $T = e^{t/2} \ge T_0$ ,  $\alpha \in [-\frac{1}{2}, \frac{1}{2}]$ :

$$\mathcal{N}_{T}(I,\alpha) \leqslant \mathcal{N}\left(\tilde{n}(2\alpha)\Phi^{t},\mathfrak{C}(I)\right) + \mathcal{N}\left(\tilde{n}(-2\alpha)\Phi^{t},\mathfrak{C}(I)\right)$$
(3.5)

and, for  $-\frac{1}{3}T^{-1/2} \leq \alpha \leq \frac{1}{3}T^{-1/2}$ ,  $\mathcal{N}_T(I,\alpha) = 0.$  (3.6)

*Proof.* The bound (3.5) follows from the more precise estimates in [2]; cf. also [9, Sect. 4]. The second statement (3.6) follows from the observation that the distance of  $\sqrt{n}$  to the nearest integer, with  $n \leq T$  and not a perfect square, is at least  $\frac{1}{2}(n+1)^{-1/2} \geq \frac{1}{2}(T+1)^{-1/2}$ .

A convenient parametrization of  $M \in G$  is given by the Iwasawa decomposition

$$M = n(u)a(v)k(\varphi) = \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix} \begin{pmatrix} v^{1/2} & 0 \\ 0 & v^{-1/2} \end{pmatrix} \begin{pmatrix} \cos\varphi & -\sin\varphi \\ \sin\varphi & \cos\varphi \end{pmatrix}$$
(3.7)

where  $\tau = u + iv$  is in the complex upper half plane  $\mathbb{H} = \{u + iv \in \mathbb{C} : v > 0\}$  and  $\varphi \in [0, 2\pi)$ . A convenient parametrization of  $g \in G'$  is then given by  $\mathbb{H} \times [0, 2\pi) \times \mathbb{R}^2$  via the decomposition

$$g = (1, \boldsymbol{\xi})n(u)a(v)k(\varphi) \eqqcolon (\tau, \varphi; \boldsymbol{\xi}).$$
(3.8)

In these coordinates, left multiplication on G becomes the group action

$$g \cdot (\tau, \varphi; \boldsymbol{\xi}) = (g\tau, \varphi_g; \boldsymbol{\xi}g^{-1})$$
(3.9)

where for

$$g = (1, \boldsymbol{m}) \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$
(3.10)

we have:

$$g\tau = u_g + iv_g = \frac{a\tau + b}{c\tau + d}$$
(3.11)

and thus

$$v_g = \text{Im}(g\tau) = \frac{v}{|c\tau + d|^2};$$
 (3.12)

furthermore

$$\varphi_g = \varphi + \arg(c\tau + d), \tag{3.13}$$

and

$$\boldsymbol{\xi}g^{-1} = (d\xi_1 - c\xi_2, -b\xi_1 + a\xi_2) - \boldsymbol{m}.$$
(3.14)

We define the abelian subgroups

$$\Gamma_{\infty} = \left\{ \begin{pmatrix} 1 & m \\ 0 & 1 \end{pmatrix} : m \in \mathbb{Z} \right\} \subset \Gamma$$

and

$$\Gamma_{\infty}' = \left\{ \left( \begin{pmatrix} 1 & m_1 \\ 0 & 1 \end{pmatrix}, (0, m_2) \right) : (m_1, m_2) \in \mathbb{Z}^2 \right\} \subset \Gamma'.$$

These subgroups are the stabilizers of the cusp at  $\infty$  of  $\Gamma \backslash G$  and  $\Gamma' \backslash G'$ , respectively.

Denote by  $\chi_R$  the characteristic function of  $[R, \infty)$  for some  $R \ge 1$ , i.e.  $\chi_R(v) = 0$  if v < R and  $\chi_R(v) = 1$  if  $v \ge R$ . For a fixed real number  $\beta$  and a continuous function  $f : \mathbb{R} \to \mathbb{R}$  of rapid decay at  $\pm \infty$ , define the function  $F_{R,\beta} : \mathbb{H} \times \mathbb{R}^2 \to \mathbb{R}$  by

$$F_{R,\beta}(\tau; \boldsymbol{\xi}) = \sum_{\gamma \in \Gamma_{\infty} \setminus \Gamma} \sum_{m \in \mathbb{Z}} f(((\boldsymbol{\xi}\gamma^{-1})_{1} + m)v_{\gamma}^{1/2})v_{\gamma}^{\beta}\chi_{R}(v_{\gamma})$$
  
$$= \sum_{\gamma \in \Gamma_{\infty}' \setminus \Gamma'} f_{\beta}(\gamma g), \qquad (3.15)$$

where  $f_{\beta}: G' \to \mathbb{R}$  is defined by

$$f_{\beta}((1,\boldsymbol{\xi})n(u)a(v)k(\varphi)) := f(\xi_1 v^{1/2})v^{\beta}\chi_R(v).$$
(3.16)

We view  $F_{R,\beta}(\tau; \boldsymbol{\xi}) = F_{R,\beta}(g)$  as a function on  $\Gamma' \backslash G'$  via the identification (3.8).

We show in [1, Sect. 3] that there is a choice of a continuous function  $f \ge 0$  with compact support, such that for  $\beta = \frac{1}{2}\sigma$ , and  $v \ge R$  with R sufficiently large, we have

$$\mathcal{N}(g,\mathfrak{C}(I))^{\sigma} \leqslant F_{R,\beta}(g) = F_{R,\beta}\left(\tau;\boldsymbol{\xi}\right). \tag{3.17}$$

The following proposition establishes under which conditions there is no escape of mass in the equidistribution of translates of non-linear horocycles. In view of Lemma 4 and (3.17), it implies (2.5) and thus Theorem 3. (Use v = 1/T and note that  $\frac{\beta}{2(\beta-1)} > \frac{1}{2}$  so the choice  $\eta = \frac{1}{2}$  is always permitted.)

**Proposition 5.** Assume f is continuous and has compact support. Let  $0 \leq \beta < \frac{3}{2}$ . Then

$$\lim_{R \to \infty} \limsup_{v \to 0} \left| \int F_{R,\beta} \left( \tilde{n}(u) a(v) \right) du \right| = 0$$
(3.18)

where the range of integration is [-1,1] for  $\beta < 1$ , and  $[-1,-\theta v^{\eta}] \cup [\theta v^{\eta},1]$  for  $\beta \ge 1$  and any  $\eta \in [0,\frac{\beta}{2(\beta-1)}), \ \theta \in (0,1).$ 

The proof of this proposition is organized in three parts: the proof for  $\beta < 1$ , a key lemma, and finally the proof for  $1 \leq \beta < \frac{3}{2}$ . In the following we assume without loss of generality that f is nonnegative, even, and that  $f(rx) \leq f(x)$  for all  $r \geq 1$  and all  $x \in \mathbb{R}$ .

# 4 Proof of Proposition 5 for $\beta < 1$

(This case is almost identical to the analogous result in [1].) Since f is rapidly decaying and  $R \ge 1$ , we have

$$F_{R,\beta}\left(\tau;\boldsymbol{\xi}\right) \ll_{f} \overline{F}_{R,\beta}\left(\tau\right) \tag{4.1}$$

where

$$\overline{F}_{R,\beta}(\tau) = \sum_{\gamma \in \Gamma_{\infty} \setminus \Gamma} v_{\gamma}^{\beta} \chi_{R}(v_{\gamma}).$$

Thus

$$\int_{-1}^{1} F_{R,\beta}\left(u+\mathrm{i}v;\boldsymbol{\xi}\right) du \ll_{f} 2 \int_{0}^{1} \overline{F}_{R,\beta}\left(u+\mathrm{i}v\right) du.$$

$$(4.2)$$

The evaluation of the integral on the right hand side is well known from the theory of Eisenstein series. We have

$$\overline{F}_{R,\beta}(\tau) = v^{\beta} \chi_{R}(v) + 2 \sum_{c=1}^{\infty} \sum_{\substack{d=1\\ \gcd(c,d)=1}}^{c-1} \sum_{m \in \mathbb{Z}} \frac{v^{\beta}}{c^{2\beta} |\tau + \frac{d}{c} + m|^{2\beta}} \chi_{R}\left(\frac{v}{c^{2} |\tau + \frac{d}{c} + m|^{2}}\right).$$
(4.3)

This function is evidently periodic in  $u = \operatorname{Re} \tau$  with period one, and its zeroth Fourier coefficient is (we denote by  $\varphi$  Euler's totient function)

$$\int_{0}^{1} \overline{F}_{R,\beta} \left( u + iv \right) du = v^{\beta} \chi_{R}(v) + 2v^{1-\beta} \sum_{c=1}^{\infty} \frac{\varphi(c)}{c^{2\beta}} \int_{\mathbb{R}} \frac{1}{(t^{2}+1)^{\beta}} \chi_{R} \left( \frac{1}{vc^{2}(t^{2}+1)} \right) dt. \quad (4.4)$$

The first term vanishes for v < R, and the second term is bounded from above by

$$2v^{1-\beta}\sum_{c=1}^{\infty}\frac{1}{c^{2\beta-1}}\int_{\mathbb{R}}\frac{1}{(t^2+1)^{\beta}}\chi_R\left(\frac{1}{vc^2(t^2+1)}\right)\,dt = 2v^{1/2}\sum_{c=1}^{\infty}K_R(cv^{1/2})\tag{4.5}$$

with the function  $K_R : \mathbb{R}_{>0} \to \mathbb{R}_{\geq 0}$  defined by

$$K_R(x) = \frac{1}{x^{2\beta - 1}} \int_{\mathbb{R}} \frac{1}{(t^2 + 1)^\beta} \chi_R\left(\frac{1}{x^2(t^2 + 1)}\right) dt.$$
(4.6)

We have  $K_R(x) \ll \max\{1, x^{2\beta-1}\} \leqslant \max\{1, x^{-1/2}\}$  and furthermore  $K_R(x) = 0$  if  $x > R^{-1/2}$ . Thus

$$\lim_{v \to 0} v^{1/2} \sum_{c=1}^{\infty} K_R(cv^{1/2}) = \int_{\mathbb{R}} K_R(x) dx, \qquad (4.7)$$

which evaluates to a constant times  $R^{-(1-\beta)}$ .

# 5 Key Lemma

**Lemma 6.** Let  $f \in C(\mathbb{R})$  be rapidly decreasing and let

$$S = \sum_{\substack{D \le c \le 2D \\ 1 \le d \le D \\ \gcd(c,d) = 1}} \sum_{m \in \mathbb{Z}} f\left(T\left(\frac{d^2}{4c} + m\right)\right).$$
(5.1)

Then, for  $D \ge 1, T > 1$  and any  $\varepsilon > 0$ , we have

$$S \ll \frac{D^2}{T^{1-\varepsilon}},\tag{5.2}$$

where the implied constant depends only on  $\varepsilon$  and f.

*Proof.* We assume without loss of generality that f is even, non-negative, and of Schwartz class. We prove two statements about S from which the statement of the Lemma will follow. They are

$$S \ll \frac{D^{2+\varepsilon'}}{T}$$
 for any  $\varepsilon' > 0$  (5.3)

and

$$S \ll \frac{D^2}{T} + D^{3/2} T^{\varepsilon''}$$
 for any  $\varepsilon'' > 0.$  (5.4)

Then S is bounded by the smaller of these expressions, and it is easy to see that the bound in (5.2) holds no matter which realizes the minimum.

Equation (5.3) is verified by summing over quadratic residues modulo 4c. Note that the conditions  $1 \leq d \leq c$  and gcd(c, d) = 1 imply  $\frac{d^2}{4c} \notin \mathbb{Z}$ . For coprime  $D \leq c \leq 2D$  and  $1 \leq d \leq D$  and  $m \in \mathbb{Z}$  such that  $\left|\frac{d^2}{4c} + m\right| \geq \frac{1}{2}$ , we use rapid decay of f to get

$$\sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant D \\ \gcd(c,d)=1}} \sum_{\substack{m \in \mathbb{Z} \\ \frac{d^2}{4c} + m \\ \frac{d^2}{2}}} f\left(T\left(\frac{d^2}{4c} + m\right)\right) \ll \sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant D \\ \frac{d^2}{4c} + m \\ \frac{d^2}{4c} + m \\ \frac{d^2}{2}}} \sum_{\substack{m \in \mathbb{Z} \\ \frac{d^2}{4c} + m \\ \frac{d$$

$$\ll \frac{1}{T^A} \sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant D}} \sum_{m \in \mathbb{Z} \setminus \{0\}} \frac{1}{|m|^A}$$
(5.6)

$$\ll \frac{D^2}{T^A} \tag{5.7}$$

for every A > 1.

For a positive integer n, denote by  $\omega(n)$  the number of distinct prime factors of n and by  $\tau(n)$  the number of divisors of n.

For coprime  $D \leq c \leq 2D$  and  $1 \leq d \leq D$  and  $m \in \mathbb{Z}$  such that  $\left|\frac{d^2}{4c} + m\right| < \frac{1}{2}$ , we have (denote by  $\|\cdot\|$  the distance to the nearest integer)

$$\sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant D \\ \gcd(c,d)=1}} \sum_{\substack{m \in \mathbb{Z} \\ |\frac{d^2}{4c} + m| < \frac{1}{2}}} f\left(T\left(\frac{d^2}{4c} + m\right)\right) \leqslant \sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant D \\ \gcd(c,d)=1}} f\left(T\left\|\frac{d^2}{4c}\right\|\right)$$
(5.8)

$$\ll \sum_{\substack{D \leqslant c \leqslant 2D\\1 \leqslant j \leqslant 2c}} 2^{\omega(4c)} f\left(T\frac{j}{4c}\right), \qquad (5.9)$$

j is a square mod 4c

since

$$#\{d \mod 4c : d^2 \equiv j \mod 4c, \ \gcd(c,d) = 1\} \ll 2^{\omega(4c)}$$
(5.10)

for every j. Now  $2^{\omega(4c)}$  is the number of squarefree divisors of 4c and is therefore at most  $\tau(4c)$ . Combined with rapid decay (we use  $f(t) \ll \frac{1}{t}$ ), this yields

$$\ll \frac{1}{T} \sum_{\substack{D \leqslant c \leqslant 2D \\ 1 \leqslant j \leqslant 4D}} \tau(4c) \frac{c}{j}.$$
(5.11)

Finally the fact that for every  $\varepsilon'' > 0$ ,  $\tau(n) \ll n^{\varepsilon''}$  gives

$$\ll \frac{D^{2+\varepsilon''} \log D}{T}$$

$$\ll \frac{D^{2+\varepsilon'}}{T}$$
(5.12)

for every  $\varepsilon' > 0$ . It suffices to note that (5.7)  $\ll$  (5.12) to verify (5.3).

Inequality (5.4) is obtained as follows. The Poisson summation formula yields

$$S \ll c_0 D^2 + \left| \sum_{\substack{n \neq 0 \\ D \leqslant c \leqslant 2D \\ 1 \leqslant d \leqslant 4c}} c_n e\left(\frac{nd^2}{4c}\right) \right|.$$
(5.13)

Here

which is

$$c_n = \frac{1}{T}\hat{f}\left(\frac{n}{T}\right) \tag{5.14}$$

where  $\hat{f}$  is the Fourier transform of f, which is also of Schwartz class, and  $e(z) = e^{2\pi i z}$  is the usual shorthand. The second term in (5.13) is bounded by

$$\left|\sum_{\substack{n\neq0\\D\leqslant c\leqslant 2D\\1\leqslant d\leqslant 4c}} c_n e\left(\frac{nd^2}{4c}\right)\right| \leqslant \sum_{r=1}^{8D} \sum_{\substack{n\neq0\\p\in d}} \sum_{\substack{4D/r\leqslant c\leqslant 8D/r\\\gcd(n,4c)=1}} \left|c_{nr}r\sum_{\substack{d \bmod 4c}} e\left(\frac{nd^2}{4c}\right)\right|$$
(5.15)

$$\leq \sum_{n \neq 0} \sum_{1 \leq r \leq 8D} \sum_{4D/r \leq c \leq 8D/r} |c_{nr}| r \sqrt{8c}.$$
(5.16)

The last inequality follows from the well known evaluation the classical Gauss sum with gcd(n, 4c) = 1

$$\sum_{d \bmod 4c} e\left(\frac{nd^2}{4c}\right) = (1+i) \varepsilon_n^{-1} \left(\frac{4c}{n}\right) \sqrt{4c},$$
(5.17)

where  $\left(\frac{4c}{n}\right)$  is the Jacobi symbol and  $\varepsilon_n = 1$  or i if n = 1 or  $3 \mod 4$ , respectively.

When |nr| < T we use the fact that the Fourier transform of f is bounded:

$$|c_{nr}| \ll \frac{1}{T}.\tag{5.18}$$

Therefore,

$$(5.16) \ll \frac{D^{3/2}}{T} \sum_{\substack{n \neq 0 \\ 1 \leqslant r \leqslant 8D \\ |nr| < T}} \frac{1}{\sqrt{r}} = \frac{D^{3/2}}{T} \sum_{1 \leqslant r \leqslant 8D} \frac{1}{\sqrt{r}} \sum_{|n| < T/r} 1 \ll D^{3/2}$$
(5.19)

When |nr| > T, we use the fact that the Fourier transform of f decays faster than any polynomial since f is smooth:

$$|c_{nr}| \ll \frac{1}{|n|^A r^A} T^{A-1}.$$
 (5.20)

We take A > 1. Then we have

$$(5.16) \ll D^{3/2} T^{A-1} \sum_{\substack{|nr|>T\\n\neq 0\\1\leqslant r\leqslant 8D}} \frac{1}{|n|^A r^{A+1/2}} \ll D^{3/2} T^{A-1} \sum_{\substack{nr>T\\n,r\geqslant 1}} \frac{1}{(nr)^A}$$
(5.21)

$$\ll D^{3/2}T^{A-1}\sum_{k=T}^{\infty} k^{-A+\varepsilon''} \ll D^{3/2}T^{\varepsilon''}.$$
 (5.22)

This proves (5.4) and the Lemma.

# 6 Proof of Proposition 5 for $\beta \ge 1$

We have

$$\tilde{n}(u)a(v) = (u + iv, 0; \boldsymbol{\xi}) \tag{6.1}$$

where  $\boldsymbol{\xi} = (u/2, -u^2/4)$ . For this choice we have

$$F_{R,\beta}(\tau; \boldsymbol{\xi}) = 2 \sum_{m \in \mathbb{Z}} f\left( (m + u^2/4) \frac{v^{1/2}}{|\tau|} \right) \frac{v^{\beta}}{|\tau|^{2\beta}} \chi_R\left(\frac{v}{|\tau|^2}\right) +$$

$$+ 2 \sum_{\substack{(c,d) \in \mathbb{Z}^2 \\ \gcd(c,d)=1 \\ c>0, d\neq 0}} \sum_{m \in \mathbb{Z}} f\left( (cu^2/4 + du/2 + m) \frac{v^{1/2}}{|c\tau+d|} \right) \frac{v^{\beta}}{|c\tau+d|^{2\beta}} \chi_R\left(\frac{v}{|c\tau+d|^2}\right).$$
(6.2)
(6.3)

The integral of the first term tends to zero as  $v \to 0$ . We write  $T = \frac{v^{1/2}}{|\tau|} = \frac{v^{1/2}}{\sqrt{u^2 + v^2}}$ . Indeed, for  $m \neq 0$  we have  $|f((m + u^2/4)T)| \ll (|m|T)^{-A}$  from rapid decay, so that

$$\int_{J} (6.2) du \ll \int_{-1}^{1} T^{-A+\beta} du \ll \int_{-1}^{1} \left(\frac{v^{1/2}}{v}\right)^{-A+\beta} du \ll v^{\frac{A-\beta}{2}}, \tag{6.4}$$

where  $J = [-1, -\theta v^{\eta}] \cup [\theta v^{\eta}, 1]$ . If  $A > \beta$ , then this contribution is negligible as  $v \to 0$ . For m = 0, we have  $|f| \ll 1$  so that the contribution of this term is, assuming  $\beta > 1$ ,

$$\ll \int_{J} T^{\beta} du \ll \int_{\theta v^{\eta}}^{1} u^{-\beta} v^{\beta/2} du = \frac{\theta^{1-\beta}}{\beta - 1} v^{\beta/2 + \eta - \beta\eta} \to 0$$
(6.5)

since  $\eta < \frac{\beta}{2(\beta-1)}$ ; similarly for  $\beta = 1$ .

It remains to analyze the contribution of (6.3). Notice that the this term is nonzero only when -d/c is in the range of integration for u, which is contained in the interval [-1, 1]. Therefore we restrict the summation to  $0 < |d| \leq c$ . Now we perform the substitution  $t = (u + d/c) v^{-1}$  to "zoom in" on each rational point and extend the range of integration to all of  $\mathbb{R}$ . This gives

$$\int_{t=-\infty}^{\infty} f\left(\left(\frac{c}{4}\left(-\frac{d}{c}+tv\right)^{2}+\frac{d}{2}\left(-\frac{d}{c}+tv\right)+m\right)\frac{1}{\sqrt{c^{2}v(t^{2}+1)}}\right)\times \frac{v}{(c^{2}v(t^{2}+1))^{\beta}}\chi_{R}\left(\frac{1}{c^{2}v(t^{2}+1)}\right)dt \quad (6.6)$$

and we need to bound

$$\sum_{c=1}^{\infty} \int_{t \in \mathbb{R}} \sum_{\substack{0 < |d| \leq c \ m \in \mathbb{Z} \\ (c,d)=1}} f\left( \left( -\frac{d^2}{4c} + m + O(ctv) \right) \frac{1}{\sqrt{c^2 v(t^2+1)}} \right) \frac{v \, dt}{(c^2 v(t^2+1))^{\beta}} \chi_R\left( \frac{1}{c^2 v(t^2+1)} \right).$$
(6.7)

Now we decompose the region  $\frac{1}{\sqrt{c^2 v(t^2+1)}} \ge \sqrt{R}$  into dyadic regions

$$2^{j} \leqslant \frac{1}{\sqrt{c^{2}v(t^{2}+1)}} < 2^{j+1}$$

for  $j \gg \log R$ . We can thus bound (6.7) by

$$\begin{split} \sum_{j\gg\log R} \sum_{c\geqslant 1} \int_{\mathbb{R}} \sum_{\substack{0<|d|\leqslant c \ m\in\mathbb{Z} \\ (c,d)=1}} f\left( \left( -\frac{d^2}{4c} + m + O(ctv) \right) \frac{1}{\sqrt{c^2 v(t^2+1)}} \right) \times \\ \times \frac{v}{(c^2 v(t^2+1))^\beta} \chi_{[2^j,2^{j+1})} \left( \frac{1}{\sqrt{c^2 v(t^2+1)}} \right) dt \end{split}$$

$$\leq \sum_{j \gg \log R} \sum_{c \geqslant 1} \int_{\mathbb{R}} \sum_{\substack{0 < |d| \leq c \ m \in \mathbb{Z}}} f\left(2^{j}\left(-\frac{d^{2}}{4c} + m + O(ctv)\right)\right) \times \\ \times \frac{v}{(c^{2}v(t^{2}+1))^{\beta}} \chi_{[2^{j},2^{j+1})}\left(\frac{1}{\sqrt{c^{2}v(t^{2}+1)}}\right) dt \\ \ll v \sum_{j \gg \log R} 2^{2\beta j} \int_{\mathbb{R}} \sum_{\substack{\frac{2^{-(j+1)}}{\sqrt{v(t^{2}+1)}} \leq c \leq \frac{2^{-j}}{\sqrt{v(t^{2}+1)}}} \sum_{\substack{0 < |d| \leq \frac{2^{-j}}{\sqrt{v(t^{2}+1)}}} \sum_{m \in \mathbb{Z}} f\left(2^{j}\left(-\frac{d^{2}}{4c} + m + O(ctv)\right)\right) dt. \\ (6.8)$$

It remains to remove the error term O(ctv) from the argument of f to apply Lemma 6. We have that  $2^j ctv \ll v^{1/2}$  for every j. Define

$$f^*(x) = \max_{-1 \le y \le 1} f(x+y).$$

Then,  $f^*(x) \ge f(x + O(v^{1/2}))$  for v small enough, and we can bound (6.8) by a similar expression with  $f^*\left(2^j\left(-\frac{d^2}{4c}+m\right)\right)$  in place of  $f\left(2^j\left(-\frac{d^2}{4c}+m+O(ctv)\right)\right)$ . To bound this we apply Lemma 6 with  $D \sim \frac{2^{-(j+1)}}{\sqrt{v(t^2+1)}}, T = 2^j$ , and  $\varepsilon = \frac{3}{2} - \beta > 0$ . Then we have

$$(6.8) \ll v \sum_{j \gg \log R} 2^{2\beta j} \int_{t \in \mathbb{R}} \frac{D^2}{T^{1-\varepsilon}} dt$$
(6.9)

$$\ll \sum_{j \gg \log R} 2^{j(2\beta - 3 + \varepsilon)} = \sum_{j \gg \log R} 2^{j(\beta - 3/2)} \to 0$$
(6.10)

as  $R \to \infty$  by our choice of  $\varepsilon$ .

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DANIEL EL-BAZ, SCHOOL OF MATHEMATICS, UNIVERSITY OF BRISTOL, BRISTOL BS8 1TW, U.K. daniel.el-baz@brisol.ac.uk

JENS MARKLOF, SCHOOL OF MATHEMATICS, UNIVERSITY OF BRISTOL, BRISTOL BS8 1TW, U.K. j.marklof@bristol.ac.uk

ILYA VINOGRADOV, SCHOOL OF MATHEMATICS, UNIVERSITY OF BRISTOL, BRISTOL BS8 1TW, U.K. ilya.vinogradov@bristol.ac.uk