# Heat Flow out of Regions in $\mathbb{R}^{\boldsymbol{m}}$ 

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## 1. Introduction

Let $D$ be an open set in $\mathbb{R}^{m}$ and let $H=H^{*} \geqq 0$ be the operator $-\Delta_{D}$ on $L^{2}(D)$ subject to Dirichlet boundary conditions, and defined by the method of quadratic forms [8, 9], so that $\operatorname{Quad}(H)=\operatorname{Dom}\left(H^{1 / 2}\right)=W_{0}^{1,2}(D)$. The heat kernel $p_{D}(x, y ; t)$ of $e^{-H t}$ is a positive $C^{\infty}$ function on $(0, \infty) \times D \times D$. It is well-known that if $D \subset E, E$ open, then

$$
\begin{equation*}
0 \leqq p_{D}(x, y ; t) \leqq p_{E}(x, y ; t) \tag{1.1}
\end{equation*}
$$

In particular putting $E=\mathbb{R}^{m}$ yields

$$
\begin{equation*}
0 \leqq p_{D}(x, y ; t) \leqq(4 \pi t)^{-m / 2} e^{-|x-y|^{2} /(4 t)} \tag{1.2}
\end{equation*}
$$

In previous papers [1, 2, 7] we have been interested in finding necessary and sufficient conditions for

$$
\begin{equation*}
Z_{D}(t)=\operatorname{trace}\left(e^{-H t}\right)=\int_{D} d x p_{D}(x, x ; t) \tag{1.3}
\end{equation*}
$$

to be finite for some $t>0$. While $Z_{D}(t)$ is always finite for $t>0$ if the volume $|D|$ of $D$ is finite this is not necessarily the case for regions with infinite volume. One reason for the importance of the function $Z_{D}(t)$ is the fact that if $Z_{D}(t)<\infty$ for $t>0$ and the asymptotic behaviour of $Z_{D}(t)$ as $t \downarrow 0$ is known, then the spectrum of $H$ is discrete: $\lambda_{1}<\lambda_{2} \leqq \lambda_{3} \leqq \ldots \leqq \lambda_{j} \leqq \ldots$ and its asymptotic distribution ( $j \uparrow \infty$ ) can be obtained via Karamata's tauberian theorem. See [1, 2, 13, 17, $18,21,22,23]$.

In this paper we will be interested in necessary and sufficient conditions on $D$ for a closely related function $Q_{D}(t)$ to be finite for some $t \geqq 0$. Here $Q_{D}(t)$ is defined by

$$
\begin{equation*}
Q_{D}(t)=\int_{D} d x \int_{D} d y p_{D}(x, y ; t) . \tag{1.4}
\end{equation*}
$$

For regions $D$ with finite volume one has in view of (1.2)

$$
\begin{equation*}
Q_{D}(t) \leqq \int_{D} d x \int_{\mathbb{R}^{m}} d y p_{D}(x, y ; t) \leqq|D|, \quad t \geqq 0 \tag{1.5}
\end{equation*}
$$

$Q_{D}(t)$ represents the amount of heat contained in $D$ at time $t$ when $D$ has temperature 1 at $t=0$ and the boundary $\partial D$ of $D$ is kept at temperature 0 for all $t>0$.

This paper is organized as follows. In Sect. 2 we will prove some elementary estimates for $Q_{D}(t)$ and $Z_{D}(t)$. In Sect. 3 and 4 we prove supercontractive estimates and gaussian upper bounds for the heat kernel $p_{D}(x, y ; t)$ under conditions on $D$ which imply compactness of the resolvent of $H$. These bounds on $p_{D}(x, y ; t)$ imply bounds on $Q_{D}(t)$. All the results which we obtain support the following.

Conjecture. If $0 \leqq t<\infty$ and $D$ is open in $\mathbb{R}^{m}$ then the following are equivalent.
(i) $Q_{D}(s)<\infty$ for all $s>t$,
(ii) $Z_{D}(s)<\infty$ for all $s>t$.

If $D$ is a horn-shaped region in $\mathbb{R}^{2}$ we prove that the conjecture holds (Theorem 5.11) and that one cannot sharpen it to the case where $s=t$ (Theorem 5.5). Finally in Sect. 6 we obtain the first two terms in the asymptotic expansion of $Q_{D}(t)$ as $t \downarrow 0$ for bounded regions $D$ in $\mathbb{R}^{m}$ with a smooth boundary. (See [3,5] for the asymptotic behaviour in some special cases).

The techniques rely on the representation (see [19]) of $\int_{D} d y p_{D}(x, y ; t)$ as a Wiener probability $P_{x}\left[T_{D}>t\right]$ that a brownian motion $x(\cdot)$ with $x(0)=x$ does not leave $D$ until $t$ :

$$
\begin{equation*}
P_{x}\left[T_{D}>t\right]=\int_{D} d y p_{D}(x, y ; t) \tag{1.6}
\end{equation*}
$$

A crucial ingredient of our calculations is the quadratic form inequality

$$
\begin{equation*}
m \int_{D} \frac{f^{2}(x)}{4 m^{2}(x)} d x \leqq \int_{D}|\nabla f(x)|^{2} d x \tag{1.7}
\end{equation*}
$$

for all $f \in C_{c}^{\infty}(D)$. Here the mean distance function $m(x)$ is defined on $D$ by

$$
\begin{equation*}
\frac{1}{m^{2}(x)}=\int_{\|u\|=1} \frac{d S(u)}{d_{u}^{2}(x)} \tag{1.8}
\end{equation*}
$$

where $d S$ is the normalized surface measure on the unit sphere of $\mathbb{R}^{m}$ and

$$
\begin{equation*}
d_{u}(x)=\min \{|t|: t \in \mathbb{R}, x+t u \notin D\}, \tag{1.9}
\end{equation*}
$$

with $d_{u}(x)=+\infty$ if the set is empty. The bound (1.7) may be found in [6, 7, 9] where it is also shown that $m(x) \geqq d(x)=\min \{|x-y|: y \notin D\}$. We say $D$ is regular if there exists a constant $c \geqq 1$ such that $d(x) \leqq m(x) \leqq c d(x)$ for all $x \in D$, and refer to $[6,7,9]$ for conditions on $\partial D$ which imply regularity. We also note that if $D$ is regular, then $H$ has compact resolvent if and only if $d(x) \rightarrow 0$
as $|x| \rightarrow \infty$. ( $H$ has compact resolvent if and only if the embedding of $W_{0}^{1,2}(D)$ into $L^{2}(D)$ is compact. More complicated necessary and sufficient conditions for this, which do not require $D$ to be regular may be found in $[12,15])$.

## 2. Elementary Estimates

Lemma 2.1. For $D$ open in $\mathbb{R}^{m}$

$$
\begin{equation*}
p_{D}(x, x ; t) \geqq(4 \pi t)^{-m / 2} e^{-m^{2} \pi^{2} t /\left(4 d^{2}(x)\right)} \tag{2.1}
\end{equation*}
$$

Proof. A cube $C_{x}$ with centre $x$ and edge length $2 d(x) / m^{1 / 2}$ is contained in $D$. Since the heat kernel for a cube is the product of $m$ one-dimensional heat kernels we obtain (2.1) by using Lemma 8 of [4].

Lemma 2.2. For $D$ open in $\mathbb{R}^{m}$

$$
\begin{equation*}
\int_{D} d y p_{D}(x, y ; t) \geqq 2^{-m} e^{-m^{2} \pi^{2} t /\left(4 d^{2}(x)\right)} \tag{2.2}
\end{equation*}
$$

Proof. By positivity and monotonicity of heat kernels

$$
\int_{D} d y p_{D}(x, y ; t) \geqq \int_{C_{x}} d y p_{D}(x, y ; t) \geqq \int_{C_{x}} d y p_{C_{x}}(x, y ; t) .
$$

By the eigenfunction expansion of $p_{C_{x}}(x, y ; t)$ :

$$
\begin{align*}
\int_{C_{x}} d y p_{C_{x}}(x, y ; t) & =\left\{\sum_{j=0}^{\infty} e^{-m t \pi^{2}(2 j+1)^{2} /\left(4 d^{2}(x)\right)} \frac{4}{\pi(2 j+1)}(-1)^{j}\right\}^{m} \\
& \geqq\left\{e^{-m t \pi^{2} /\left(4 d^{2}(x)\right)} \frac{4}{\pi}-e^{-m t \pi^{2} 9 /\left(4 d^{2}(x)\right)} \frac{4}{3 \pi}\right\}^{m} \\
& \geqq 2^{-m} e^{-m^{2} t \pi^{2} /\left(4 d^{2}(x)\right)} . \tag{2.3}
\end{align*}
$$

By combining the above lemmas we obtain the following.
Corollary 2.3. If either $Z_{D}(t)<\infty$ or $Q_{D}(t)<\infty$ then

$$
\begin{equation*}
\int_{D} d x e^{-m^{2} \pi^{2} t /\left(4 d^{2}(x)\right)}<\infty \tag{2.4}
\end{equation*}
$$

The following lemma gives an upper bound on $Q_{D}(t)$ in terms of $p_{D}(x, x ; t)$ and the normalized eigenfunction in $L^{2}(D)$ corresponding to $\lambda_{1}$.

Lemma 2.4. Suppose the spectrum of $-\Delta_{D}$ is discrete and $\phi_{1}$ is the normalized eigenfunction in $L^{2}(D)$ corresponding to $\lambda_{1}$. Then

$$
\begin{equation*}
Q_{D}(t) \leqq e^{t \lambda_{1}}\left\{\int_{D} p_{D}(x, x ; t)\left(\phi_{1}(x)\right)^{-1} d x\right\}^{2} \tag{2.5}
\end{equation*}
$$

Proof. Since $p_{D}(x, y ; t)$ is of positive type

$$
\begin{equation*}
p_{D}(x, y ; t) \leqq\left(p_{D}(x, x ; t) p_{D}(y, y ; t)\right)^{1 / 2} \tag{2.6}
\end{equation*}
$$

Hence

$$
\begin{equation*}
Q_{D}(t) \leqq\left\{\int_{D} d x\left(p_{D}(x, x ; t)\right)^{1 / 2}\right\}^{2} \tag{2.7}
\end{equation*}
$$

Using the eigenfunction expansion and the positivity of $\phi_{1}$

$$
\begin{align*}
\left(p_{D}(x, x ; t)\right)^{1 / 2} & =\left\{\sum_{j=1}^{\infty} e^{-t \lambda_{j}}\left(\phi_{j}(x)\right)^{2}\right\}^{1 / 2} \\
& \geqq e^{-t \lambda_{1} / 2} \phi_{1}(x) \tag{2.8}
\end{align*}
$$

so that

$$
\begin{equation*}
Q_{D}(t) \leqq\left\{\int_{D} d x p_{D}(x, x ; t) e^{I \lambda_{1} / 2}\left(\phi_{1}(x)\right)^{-1}\right\}^{2} \tag{2.9}
\end{equation*}
$$

The following proposition was proved in [7]. However, the first part follows more directly from Lemma 2.1.

Proposition 2.5. If $D$ is an arbitrary open set in $\mathbb{R}^{m}$ then the first of the following conditions implies the second. If $D$ is regular then the two conditions are equivalent.
(i) $Z_{D}(t)<\infty$ for all $t>0$,
(ii) $\int_{D} e^{-t / d^{2}(x)} d x<\infty$ for all $t>0$.

Much of our analysis is motivated by the attempt to find something close to a converse of the following result.

Lemma 2.6. If $D$ is open in $\mathbb{R}^{m}$ and $Q_{D}(t)<\infty$ for some $t>0$, then $Z_{D}(s)<\infty$ for all $s>t$.

Proof. If $\varepsilon>0$ then

$$
\begin{align*}
Z_{D}(t+\varepsilon) & =\int_{D} p_{D}(x, x ; t+\varepsilon) d x \\
& =\int_{D} d x \int_{D} d y p_{D}(x, y ; \varepsilon) p_{D}(y, x ; t) \\
& \leqq(4 \pi \varepsilon)^{-m / 2} Q_{D}(t) \tag{2.10}
\end{align*}
$$

## 3. Supercontractive Estimates

Throughout this section we shall take $\varphi$ to be the function

$$
\begin{equation*}
\varphi(x)=\left(1+x^{2}\right)^{-\alpha} \tag{3.1}
\end{equation*}
$$

on an open set $D$ in $\mathbb{R}^{m}$, where $\alpha>m / 2$. Our key assumption in this section, that

$$
\begin{equation*}
m(x)^{2}=o\left(\frac{1}{\log \left(1+x^{2}\right)}\right) \tag{3.2}
\end{equation*}
$$

as $|x| \rightarrow \infty, x \in D$, should be compared with (2.4). Lemma 3.1 and Theorem 3.4 are fairly close to being converses to Corollary 2.3.
Lemma 3.1. If $D$ is open and regular and (3.2) holds then $Z_{D}(t)<\infty$ for all $t>0$.
Proof. Since $m(x)$ and $d(x)$ are of the same order of magnitude, (3.2) implies that for any $u>0$ there exists $v>0$ such that

$$
\begin{equation*}
\frac{1}{d^{2}(x)} \geqq u \log \left(1+x^{2}\right)-v, \quad x \in D . \tag{3.3}
\end{equation*}
$$

Therefore

$$
\begin{equation*}
\int_{D} e^{-t / d^{2}(x)} d x \leqq \int_{D} e^{-t u \log \left(1+x^{2}\right)+t v} d x \leqq e^{t v} \int_{\mathbb{R}^{m}}\left(1+x^{2}\right)^{-t u} d x<\infty \tag{3.4}
\end{equation*}
$$

provided $u>m /(2 t)$. The proof is completed by applying Proposition 2.5.
Lemma 3.2. The function $\varphi$ defined in (3.1) satisfies

$$
\begin{gather*}
\int_{D} \varphi(x) d x<\infty  \tag{3.5}\\
|\Delta \varphi| \leqq c_{1} \varphi \tag{3.6}
\end{gather*}
$$

for some $c_{1}<\infty$ and all $x \in D$. Furthermore for all $\varepsilon>0$ there exists $\gamma>0$ such that for all $x \in D$

$$
\begin{equation*}
-\log \varphi \leqq \frac{\varepsilon}{m^{2}(x)}+\gamma \tag{3.7}
\end{equation*}
$$

Proof. Inequality (3.5) follows from the assumption that $\alpha>m / 2$. Inequality (3.6) is a direct computation using the formula

$$
\begin{equation*}
\Delta \varphi=\frac{d^{2} \varphi}{d r^{2}}+\frac{m-1}{r} \frac{d \varphi}{d r} \tag{3.8}
\end{equation*}
$$

Finally (3.7) follows immediately from (3.1) and (3.2).
We now follow the standard procedure [9] of transferring the problem to the weighted space $L^{2}\left(D, \varphi^{2} d x\right)$. If we put $V=\Delta \varphi / \varphi$ then $V$ is bounded and $(-\Delta+V) \varphi=0$. We define the unitary operator $U$ from $L^{2}\left(D, \varphi^{2} d x\right)$ to $L^{2}(D, d x)$ by $U f=\varphi f$ and consider the quadratic form $Q_{\varphi}$ defined on the subspace $C_{c}^{\infty}(D)$ of $L^{2}\left(D, \varphi^{2} d x\right)$ by

$$
\begin{equation*}
Q_{\varphi}(f)=\int_{D}|\nabla f|^{2} \varphi^{2} d x \tag{3.9}
\end{equation*}
$$

This form is closable and its closure is associated with the self-adjoint operator

$$
\begin{equation*}
0 \leqq H_{\varphi}=U^{-1}(-\Delta+V) U \tag{3.10}
\end{equation*}
$$

on $L^{2}\left(D, \varphi^{2} d x\right)$. The condition $\varphi \in L^{2}(D, d x)$ often imposed in this argument is actually irrelevant.
Theorem 3.3. The operator $e^{-H_{\varphi} t}$ is bounded from the weighted space $L^{p}\left(D, \varphi^{2} d x\right)$ to the weighted space $L^{q}\left(D, \varphi^{2} d x\right)$ for all $1<p \leqq q<\infty$ and all $0<t<\infty$.

Proof. Inequality (3.7) implies that if $\varepsilon>0$ then

$$
\begin{equation*}
-\log \varphi \leqq \varepsilon(-\Delta+V)+\mu \tag{3.11}
\end{equation*}
$$

as a quadratic form inequality on $L^{2}(D, d x)$ for some $\mu>0$. Rosen's lemma $[9,16]$ now establishes that for all $\varepsilon>0$ there exists $\beta(\varepsilon)<\infty$ such that

$$
\begin{equation*}
\int_{D}\left(f^{2} \log f\right) \varphi^{2} d x \leqq \varepsilon Q_{\varphi}(f)+\beta(\varepsilon)\|f\|_{2}^{2}+\|f\|_{2}^{2} \log \|f\|_{2} \tag{3.12}
\end{equation*}
$$

for all $0 \leqq f \in \operatorname{Quad}\left(H_{\varphi}\right) \cap L^{1} \cap L^{\infty}$. It follows by [16] or by a simplified version of Theorem 2.2.7 of [9] that $e^{-H_{\varphi p} t}$ is bounded from $L^{2}\left(D, \varphi^{2} d x\right)$ to $L^{q}\left(D, \varphi^{2} d x\right)$ for all $2 \leqq q<\infty$ and $t>0$. The result now follows by duality and the fact that $e^{-H_{\varphi} t}$ is a self-adjoint semigroup.
Theorem 3.4. For $D$ open in $\mathbb{R}^{m}$ satisfying (3.2) one has $Q_{D}(t)<\infty$ for all $t>0$.
Proof. If $p^{\prime}$ is the integral kernel of $e^{-(-\Delta+V) t}$ and $p_{\varphi}$ is the integral kernel of $e^{-H_{\varphi} t}$ then

$$
\begin{align*}
0 \leqq p_{D}(x, y ; t) & \leqq e^{t\|V\|_{\infty}} p^{\prime}(x, y ; t) \\
& =e^{t\|V\|_{\infty}} p_{\varphi}(x, y ; t) \varphi(x) \varphi(y) \tag{3.13}
\end{align*}
$$

Therefore

$$
\begin{align*}
Q_{D}(t) & \leqq e^{t\|V\|_{\infty}} \int_{D} d x \int_{D} d y p_{\varphi}(x, y ; t)(\varphi(x))^{-1}(\varphi(y))^{-1}(\varphi(x))^{2}(\varphi(y))^{2} \\
& =e^{t\|V\|_{\infty}}\left\langle e^{-H_{\varphi} t} \varphi^{-1}, \varphi^{-1}\right\rangle \\
& \leqq e^{t\|V\|_{\infty}}\left\|e^{-H_{\varphi} t}\right\|_{L^{p} \rightarrow L^{q}} \cdot\left\|\varphi^{-1}\right\|_{L^{p}}^{2} \tag{3.14}
\end{align*}
$$

provided $1<p<2$ and $p^{-1}+q^{-1}=1$. But

$$
\begin{align*}
\left\|\varphi^{-1}\right\|_{p}^{p} & =\int_{D} \varphi^{-p} \varphi^{2} d x \\
& \leqq \int_{D}\left(1+x^{2}\right)^{-\alpha(2-p)} d x<\infty \tag{3.15}
\end{align*}
$$

provided $p-1>0$ is small enough.
It is clear that if we replace (3.2) by a stronger hypotheses, it will be possible to obtain sharper information about the heat kernel $p_{D}$ and in particular to
estimate the rate at which $Q_{D}(t)$ diverges as $t \downarrow 0$. One such specialization of the above calculations is given in Sect. 4, but we emphasize that many intermediate situations may also be considered.

## 4. Gaussian Upper Bounds on the Heat Kernel

In this section we obtain further upper bounds on the heat kernel under stronger hypotheses. Throughout this section we assume that

$$
\begin{equation*}
m(x)^{2} \leqq c\left(1+x^{2}\right)^{-\alpha} \tag{4.1}
\end{equation*}
$$

for some $c>0, \alpha>0$ and all $x$ in an open set $D$. We also put

$$
\begin{equation*}
\varphi(x)=\left(1+x^{2}\right)^{-\beta} \tag{4.2}
\end{equation*}
$$

for some $\beta>0$.
Lemma 4.1. For D open in $\mathbb{R}^{m}$ and assuming hypotheses (4.1), (4.2), we have

$$
\begin{equation*}
0 \leqq p_{D}(x, y ; t) \leqq c \varphi(x) \varphi(y) t^{-\gamma / 2} \tag{4.3}
\end{equation*}
$$

for all $0<t \leqq 1$ and $x, y \in D$, where

$$
\begin{equation*}
\gamma=m+4 \beta / \alpha \tag{4.4}
\end{equation*}
$$

Proof. Direct calculations lead to the bounds

$$
\begin{gather*}
|\Delta \varphi(x)| \leqq c_{3} \varphi(x),  \tag{4.5}\\
-\log \varphi(x) \leqq \frac{\varepsilon}{m(x)^{2}}+c_{4}-\frac{\beta}{\alpha} \log \varepsilon \tag{4.6}
\end{gather*}
$$

for all $x \in D$ and $0<\varepsilon<\infty$. Applying Rosen's lemma as in the proof of Theorem 3.3 we see that (3.12) holds for all $0<\varepsilon \leqq 1$, with

$$
\begin{equation*}
\beta(\varepsilon)=a_{2}-\frac{\gamma}{4} \log \varepsilon . \tag{4.7}
\end{equation*}
$$

By $[9,11]$ we deduce that

$$
\begin{equation*}
0 \leqq p_{\varphi}(x, y ; t) \leqq a_{3} t^{-\gamma / 2} \tag{4.8}
\end{equation*}
$$

for all $0<t \leqq 1$. This implies the claimed result as in the proof of Theorem 3.3.
Theorem 4.2. Let $D$ be open and (4.1), (4.2) hold, let $E \geqq 0$ be the bottom of the spectrum of $H$ acting on $L^{2}(D)$ and let $0<\delta<1$. Then

$$
\begin{equation*}
0<p_{D}(x, y ; t) \leqq c_{\delta} t^{-\gamma / 2} e^{(\delta-E) t-d^{2}(x, y) /(4 t+4 \delta t)} \varphi(x) \varphi(y), \tag{4.9}
\end{equation*}
$$

for all $x \in D, y \in D, 0<t<\infty$ where

$$
\begin{equation*}
d(x, y)=\sup \{|\psi(x)-\psi(y)|: \psi \in E\} \tag{4.10}
\end{equation*}
$$

and

$$
\begin{equation*}
E=\left\{\psi: D \rightarrow \mathbb{R} \quad \text { such that }\|\nabla \psi\|_{\infty} \leqq 1\right\} . \tag{4.11}
\end{equation*}
$$

Proof. This is a straightforward application of Theorem 4 of [10]. The replacement of the dimension of $D$ in that paper by $\gamma$ causes no difficulties. The quantity $d$ in that paper is called the riemannian distance between $x$ and $y$ in $D$, but (4.10) is the more precise definition.

For gaussian lower bounds on the heat kernel we refer to Theorem 3 of [4].

It is always the case that $d(x, y) \geqq|x-y|$. The example of the region $S_{\lambda}$ in Theorem 5.10 shows that $d(x, y)$ may be much larger than the euclidean distance. We note that in that example (4.1) holds with $\alpha=\frac{1}{\lambda}-1>0$ so that Theorem 4.2 is applicable to this region. Note also that $d(x, 0) \sim|x|^{1 / \lambda}$ as $|x| \rightarrow \infty$ since any curve from 0 to $x$ in $S_{\lambda}$ must follow the spiral.

## 5. Horn-shaped Regions

In previous papers [1, 2] we have obtained a theorem for $Z_{F}(t)$, where $F$ is a horn-shaped region in $\mathbb{R}^{m}$. In this section we prove a corresponding theorem for $Q_{F}(t)$. First we recall the notation and definitions of [1].

Notation. A point in $\mathbb{R}^{m}(m=2,3 \ldots)$ is denoted by $(x, y)$ where $y \in \mathbb{R}^{m-1}$ (orthogonal to the $x$-axis). Let $P_{x}$ be the plane through $(x, 0)$ orthogonal to the $x$ axis and let $F(x)$ be the orthogonal projection of $P_{x} \cap F$ onto $P_{0}$ where $F$ is an open set in $\mathbb{R}^{m}$.

Definition 5.1. An open set $F$ in $\mathbb{R}^{m}$ is (one-sided) horn-shaped if
(1) $F$ is connected,
(2) $F(x) \subset F\left(x^{\prime}\right)$ for all $x \geqq x^{\prime}>0, F(x)$ is empty for $x \leqq 0$,
(3) $\int_{0}^{\delta}|F(x)| d x<\infty$ for $\delta \in[0, \infty)$.
$(|F(x)|$ is the $(m-1)$-dimensional volume of $F(x))$.
Definition 5.2. Let $p_{F}\left(\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right) ; t\right)$ be the heat kernel for $-\Delta_{F}+\frac{\partial}{\partial t}$ and let $p_{F(x)}\left(y_{1}, y_{2} ; t\right)$ be the heat kernel for $-\Delta_{F(x)}+\frac{\partial}{\partial t}$, where $-\Delta_{F(x)}$ is the ( $m$
-1)-dimensional Dirichlet laplacian for $F(x), y_{1}, y_{2} \in F(x)$. Let $P_{(x, y)}\left[T_{F}>t\right]$ be the probability with respect to Wiener measure that a brownian motion $(x(\cdot)$, $y(\cdot))$ in $\mathbb{R}^{m}$ with $(x(0), y(0))=(x, y) \in F$ does not leave $F$ up to time $t$. We define $P_{y}\left[T_{F(x)}>t\right]$ for a brownian motion $y(\cdot)$ with $y(0)=y \in F(x)$ in $\mathbb{R}^{m-1}$ similarly.

By Definition 5.2, (1.4) and (1.9) we have

$$
\begin{gather*}
P_{(x, y)}\left[T_{F}>t\right]=\int_{0}^{\infty} d x_{1} \int_{F\left(x_{1}\right)} d y_{1} p_{F}\left((x, y),\left(x_{1}, y_{1}\right) ; t\right)  \tag{5.1}\\
P_{y}\left[T_{F(x)}>t\right]=\int_{F(x)} d y_{1} p_{F(x)}\left(y, y_{1} ; t\right)  \tag{5.2}\\
Q_{F(x)}(t)=\int_{F(x)} d y P_{y}\left[T_{F(x)}>t\right] \tag{5.3}
\end{gather*}
$$

Theorem 5.3. Let $F$ be (one-sided) horn-shaped in $\mathbb{R}^{m}(m=2,3, \ldots)$. Then for all $t$ for which $\int_{0}^{\infty} d x Q_{F(x)}(t)$ is finite one has

$$
\begin{equation*}
\int_{0}^{\infty} d x Q_{F(x)}(t)-\frac{2}{(\pi t)^{1 / 2}} \int_{0}^{\infty} d x Q_{F(x)}(t) \int_{x}^{\infty} d q e^{-q^{2} /(4 t)} \leqq Q_{F}(t) \leqq \int_{0}^{\infty} d x Q_{F(x)}(t) \tag{5.4}
\end{equation*}
$$

The proof of this theorem starts at Lemma 5.7. For horn-shaped regions in $\mathbb{R}^{2}$ it implies the following.

Theorem 5.4. Let $f_{1}, f_{2}$ be positive, continuous and decreasing on $(0, \infty)$ such that

$$
\begin{gather*}
\int_{0}^{\delta}\left(f_{1}(x)+f_{2}(x)\right) d x<\infty, \quad \delta \in(0, \infty),  \tag{5.5}\\
\int_{1}^{\infty} d x\left(f_{1}(x)+f_{2}(x)\right) e^{-t \pi^{2} /\left(f_{1}(x)+f_{2}(x)\right)^{2}}<\infty, \tag{5.6}
\end{gather*}
$$

and let

$$
\begin{equation*}
F=\left\{(x, y) \mid x>0,-f_{1}(x)<y<f_{2}(x)\right\} . \tag{5.7}
\end{equation*}
$$

Then

$$
\begin{equation*}
\left|Q_{F}(t)-\frac{8}{\pi^{2}} \int_{0}^{\infty} d x f(x) \sum_{k=1,3, \ldots} k^{-2} e^{-t \pi^{2} k^{2} / f^{2}(x)}\right| \leqq 4 \int_{0}^{t / 2} d x f(x), \tag{5.8}
\end{equation*}
$$

where

$$
\begin{equation*}
f(x)=f_{1}(x)+f_{2}(x) \tag{5.9}
\end{equation*}
$$

Proof. Since $f_{1}, f_{2}$ are positive, continuous and decreasing and satisfy (5.5) $H$ is horn-shaped. Furthermore $F(x)=\left(-f_{1}(x), f_{2}(x)\right)$ and $p_{F(x)}\left(y_{1}, y_{2} ; t\right)$ is the onedimensional heat kernel with eigenfunction expansion

$$
\begin{equation*}
p_{F(x)}\left(y_{1}, y_{2} ; t\right)=\frac{2}{f(x)} \sum_{k=1}^{\infty} e^{-t \pi^{2} k^{2} / f^{2}(x)}\left(\sin \frac{\pi k\left(f_{1}(x)+y_{1}\right)}{f(x)}\right)\left(\sin \frac{\pi k\left(f_{1}(x)+y_{2}\right)}{f(x)}\right) . \tag{5.10}
\end{equation*}
$$

Hence by (5.3) and (5.10)

$$
\begin{equation*}
Q_{F(x)}(t)=\frac{8 f(x)}{\pi^{2}} \sum_{k=1,3, \ldots} k^{-2} e^{-t \pi^{2} k^{2} / f^{2}(x)} \tag{5.11}
\end{equation*}
$$

Finally $Q_{F(x)} \leqq f(x)$ by (1.5) so that

$$
\begin{align*}
& \frac{2}{(\pi t)^{1 / 2}} \int_{0}^{\infty} d x Q_{F(x)}(t) \int_{x}^{\infty} d q e^{-q^{2} /(4 t)} \leqq \frac{2}{(\pi t)^{1 / 2}} \int_{0}^{\infty} d x f(x) \int_{x}^{\infty} d q e^{-q^{2} /(4 t)} \\
& \quad \leqq \frac{2}{(\pi t)^{1 / 2}} \int_{0}^{t^{1 / 2}} d x f(x) \int_{0}^{\infty} d q e^{-q^{2} /(4 t)}+\frac{2}{(\pi t)^{1 / 2}} \int_{t^{1 / 2}}^{\infty} d x f\left(t^{1 / 2}\right) \int_{x}^{\infty} d q e^{-q^{2} /(4 t)} \\
& \quad=2 \int_{0}^{t^{1 / 2}} d x f(x)+4 \pi^{-1 / 2} t^{1 / 2} f\left(t^{1 / 2}\right) \int_{1}^{\infty} d q q^{-2} e^{-q^{2} / 4} \leqq 4 \int_{0}^{t^{1 / 2}} d x f(x) \tag{5.12}
\end{align*}
$$

Theorem 5.5. Let $F$ be as in Theorem 5.4 with $f_{1}(x)=0$ and

$$
\begin{equation*}
f(x)=f_{2}(x)=\pi\{\log (3+x)+\beta \log \log (3+x)\}^{-1 / 2}, \quad \beta>0, x>0 \tag{5.13}
\end{equation*}
$$

Then

$$
\begin{equation*}
Q_{F}(t)<\infty \quad \text { if and only if } t>1, \beta>0 \text { or } t=1, \beta>\frac{1}{2} \tag{5.14}
\end{equation*}
$$

and

$$
\begin{equation*}
Z_{F}(t)<\infty \quad \text { if and only if } t>1, \beta>0 \quad \text { or } \quad t=1, \beta>1 \tag{5.15}
\end{equation*}
$$

Proof. The term $f(x) e^{-t \pi^{2} / f^{2}(x)}$ behaves like $(3+x)^{-t}(\log (3+x))^{-\beta t-1 / 2}(1+o(1))$ for $x \rightarrow+\infty$. Hence $f(x) e^{-t \pi^{2} / \delta^{2}(x)}$ is integrable on $[0, \infty)$ if and only if $t>1$, $\beta>0$ or $t=1$ and $\beta>\frac{1}{2}$. The remainder series $\sum_{k=3,5, \ldots} f(x) e^{-t \pi^{2} k^{2} / f^{2}(x)} / k^{2}$ is integrable for all $t>\frac{1}{9}$. To prove (5.15) we use Theorem 3 and (43) from [2]

$$
\begin{equation*}
\left|Z_{F}(t)-\frac{1}{(4 \pi t)^{1 / 2}} \int_{0}^{\infty} d x \sum_{k=1}^{\infty} e^{-t \pi^{2} k^{2} / f^{2}(x)}\right| \leqq \frac{f(0)}{4(\pi t)^{1 / 2}} \tag{5.16}
\end{equation*}
$$

The leading term $e^{-t \pi^{2} / f^{2}(x)}$ behaves like $(3+x)^{-t}(\log (3+x))^{-\beta t}(1+o(1))$ for $x$ $\rightarrow \infty$. Hence $e^{-i \pi^{2} / f^{2}(x)}$ is integrable on $[0, \infty)$ if and only if $t>1, \beta>0$ or $t=1$, $\beta>1$. The remaining series $\sum_{k=2}^{\infty} e^{-t \pi^{2} k^{2} / f^{2}(x)}$ is integrable for all $t>\frac{1}{4}$.

Suppose $F$ is horn-shaped as in Theorem 5.4 and $f$ is integrable on $[0, \infty)$. Then

$$
\begin{equation*}
\lim _{t \downarrow 0} Q_{F}(t)=\int_{0}^{\infty} d x f(x) \tag{5.17}
\end{equation*}
$$

In this case Theorem 5.4 provides the leading term of

$$
Q_{F}(t)-\int_{0}^{\infty} d x f(x) \quad \text { as } t \downarrow 0
$$

Example 5.6. Let $F=\left\{(x, y) \mid x>0,0<y<e^{-x}\right\}$. Then

$$
\begin{equation*}
\left|Q_{F}(t)-1-\frac{2 t^{1 / 2}}{\pi^{1 / 2}} \log t\right| \leqq 14 t^{1 / 2}, \quad 0<t \leqq 1 \tag{5.18}
\end{equation*}
$$

Proof. Note that $y \rightarrow\left(e^{-\alpha y^{2}}-1\right) / y^{2}$ is strictly increasing on $[0, \infty)$ for $\alpha>0$. By Theorems 5.3 and 5.4

$$
\begin{align*}
Q_{F}(t) & \leqq \sum_{k=1,3, \ldots} \frac{8}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x-t \pi^{2} k^{2} e^{2 x}} \\
& =1+\sum_{k=1,3, \ldots} \frac{8}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right) \\
& \leqq 1+\sum_{k=1}^{\infty} \frac{4}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right) \\
& \leqq 1+\int_{1}^{\infty} d k \frac{4}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right) \tag{5.19}
\end{align*}
$$

Integration by parts with respect to $k$ gives

$$
\begin{align*}
Q_{F}(t) & \leqq 1-8 t \int_{1}^{\infty} d k \int_{0}^{\infty} d x e^{x-t \pi^{2} k^{2} e^{2 x}} \\
& =1-8 t \int_{1}^{\infty} d x x^{-1} \int_{x}^{\infty} d k e^{-t \pi^{2} k^{2}} \\
& =1-8 t \int_{1}^{\infty} d x(\log x) e^{-t \pi^{2} x^{2}} \\
& \leqq 1-8 t \int_{0}^{\infty} d x(\log x) e^{-t \pi^{2} x^{2}} \\
& \leqq 1+(4 t / \pi)^{1 / 2}\left(C+\log \left(4 \pi^{2} t\right)\right) \tag{5.20}
\end{align*}
$$

by 4.333 of [14]. $C$ is Euler's constant. By Theorem 5.4

$$
\begin{align*}
Q_{F}(t) \geqq & \geqq+\sum_{k=1,3, \ldots .} \frac{8}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right)-4 t^{1 / 2} \\
\geqq & 1+\sum_{k=2}^{\infty} \frac{4}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right)-4 t^{1 / 2} \\
& +\frac{8}{\pi^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} e^{2 x}}-1\right) \\
\geqq & 1+\int_{1}^{\infty} d k \cdot \frac{4}{\pi^{2} k^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} k^{2} e^{2 x}}-1\right)-4 t^{1 / 2} \\
& +\frac{8}{\pi^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} e^{2 x}}-1\right) \\
= & 1+(4 t / \pi)^{1 / 2}\left(C+\log \left(4 \pi^{2} t\right)\right)-4 t^{1 / 2} \\
& +\frac{12}{\pi^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} e^{2 x}}-1\right)+8 t \int_{0}^{1} d x(\log x) e^{-t \pi^{2} x^{2}} . \tag{5.21}
\end{align*}
$$

Moreover by 3.464 of [14]

$$
\begin{aligned}
\frac{12}{\pi^{2}} \int_{0}^{\infty} d x e^{-x}\left(e^{-t \pi^{2} e^{2 x}}-1\right) & =\frac{12}{\pi^{2}} \int_{1}^{\infty} d x x^{-2}\left(e^{-t \pi^{2} x^{2}}-1\right) \\
& \geqq \frac{12}{\pi^{2}} \int_{0}^{\infty} d x x^{-2}\left(e^{-t \pi^{2} x^{2}}-1\right)=12 \pi^{-1 / 2} t^{1 / 2}
\end{aligned}
$$

and for $0<t \leqq 1$

$$
8 t \int_{0}^{1} d x(\log x) e^{-t \pi^{2} x^{2}} \geqq-8 t^{1 / 2}
$$

Combining these inequalities with (5.21) we obtain

$$
\begin{equation*}
Q_{F}(t) \geqq 1+(4 t / \pi)^{1 / 2}\left(C+\log \left(4 \pi^{2} t\right)-6 \pi^{1 / 2}-6\right), \quad 0<t \leqq 1 \tag{5.22}
\end{equation*}
$$

which completes the proof.
Lemma 5.7. Let $F$ be horn-shaped in $\mathbb{R}^{m}$ and $(x, y) \in F$. Then

$$
\begin{equation*}
P_{(x, y)}\left[T_{F}>t\right] \leqq \int_{0}^{x} d \xi \rho(\xi ; t) P_{y}\left[T_{F(x-\xi)}>t\right] \tag{5.23}
\end{equation*}
$$

where $\rho(\xi ; t)$ is given by

$$
\rho(\xi ; t)=\left\{\begin{array}{ll}
0 & \xi \leqq 0  \tag{5.24}\\
(\pi t)^{-1 / 2} e^{-\xi^{2} /(4 t)}, & \xi>0
\end{array} .\right.
$$

Proof. Let $\omega=(x(\cdot), y(\cdot))$ be a brownian path in $\mathbb{R}^{m}$ with $(x(0), y(0))=(x, y)$. Then since $F$ is horn-shaped

$$
\begin{aligned}
& \{\omega:(x(\tau), y(\tau)) \in F, 0 \leqq \tau \leqq t\} \\
& \quad \subset \bigcup_{\{\xi: 0 \leqq \xi \leqq x\}}\left\{\omega: \min _{0 \leqq \tau \leqq t} x(\tau)>x-\xi, y(\tau) \in F(x-\xi), 0 \leqq \tau \leqq t\right\}
\end{aligned}
$$

Since $x(\cdot)$ and $y(\cdot)$ are independent and the random variable $\min _{0 \leq t \leq t} x(\tau)$ has probability density $\rho(\xi ; t)$ the lemma follows.

Lemma 5.8. Let $F$ be horn-shaped in $\mathbb{R}^{m},(x, y) \in F$. Then

$$
\begin{equation*}
P_{(x, y)}\left[T_{F}>t\right] \geqq \int_{\{\xi ; \xi \geq 0, y \in F(x+\xi)\}} d \xi P_{y}\left[T_{F(x+\xi)}>t\right] \frac{\partial}{\partial \xi} P_{x}\left[T_{(0, x+\xi)}>t\right], \tag{5.25}
\end{equation*}
$$

where

$$
\begin{equation*}
P_{x}\left[T_{(0, x+\xi)}>t\right]=P_{x}[0<x(\tau)<x+\xi, 0 \leqq \tau \leqq t] . \tag{5.26}
\end{equation*}
$$

Proof. Let $\omega$ be as in the proof of Lemma 5.7. Then since $F$ is horn-shaped

$$
\begin{align*}
& \{\omega:(x(\tau), y(\tau)) \in F, 0 \leqq \tau \leqq t\} \\
& \quad \supset \bigcup_{\xi \leqq 0}\left\{\omega: \min _{0 \leqq \tau \leqq t} x(\tau) \geqq 0, \max _{0 \leqq \tau \leqq t} x(\tau) \leqq x+\xi, y(\tau) \in F(x+\xi), 0 \leqq \tau \leqq t\right\} . \tag{5.27}
\end{align*}
$$

Since $x(\cdot), y(\cdot)$ are independent and the random variable $\max _{0 \leqq \tau \leqq t} x(\tau)$ conditioned to $\min _{0 \leqq \tau \leqq t} x(\tau)>0$ has probability density $\frac{\partial}{\partial \xi} P_{x}\left[T_{(0, x+\xi)}>t\right] \begin{gathered}0 \leqq \tau \leq t \\ \text { on }[x, \infty)\end{gathered}$ the lemma follows.

Lemma 5.9. For $x>0, \xi>0$

$$
\begin{equation*}
P_{x}\left[T_{(0, x+\xi)}>t\right] \geqq 1-\int_{x}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime}-\int_{\xi}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime} . \tag{5.28}
\end{equation*}
$$

Proof.

$$
\begin{align*}
P_{x}\left[T_{(0, x+\xi)}>t\right] & =P_{x}[0<x(\tau)<x+\xi, 0 \leqq \tau \leqq t] \\
& \geqq P_{x}\left[\max _{0 \leqq \tau \leqq t} x(\tau)<x+\xi\right]-P_{x}\left[\min _{0 \leqq \tau \leqq t} x(\tau)<0\right] \\
& =\int_{0}^{\xi} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime}-\int_{x}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime} . \tag{5.29}
\end{align*}
$$

Proof of Theorem 5.3. By Lemma 5.7 and Fubini's theorem

$$
\begin{align*}
\int_{F(x)} d y P_{(x, y)}\left[T_{F}>t\right] & \leqq \int_{0}^{x} d \xi \rho(\xi ; t) \int_{F(x)} d y P_{y}\left[T_{F(x-\xi)}>t\right] \\
& \leqq \int_{0}^{x} d \xi \rho(\xi ; t) \int_{F(x-\xi)} d y P_{y}\left[T_{F(x-\xi)}>t\right] \\
& =\left\{\rho(\cdot ; t) * Q_{F(\cdot)}(t)\right\}(x) \tag{5.30}
\end{align*}
$$

where $*$ denotes convolution (with respect to $x$ ). Hence

$$
\begin{equation*}
Q_{F}(t) \leqq \int_{0}^{\infty} d x \rho(\cdot ; t) * Q_{F(\cdot)}(t)(x)=\int_{0}^{\infty} d x Q_{F(x)}(t) \tag{5.31}
\end{equation*}
$$

We define $P_{y}\left[T_{F(x+\xi)}>t\right]$ to be zero if $y \notin F(x+\xi)$. Then by Lemma 5.8, Fubini's theorem and integration by parts we have

$$
\begin{align*}
\int_{F(x)} d y P_{(x, y)}\left[T_{F}>t\right] & \geqq \int_{F(x)} d y \int_{[0, \infty)} d \xi P_{y}\left[T_{F(x+\xi)}>t\right] \frac{\partial}{\partial \xi} P_{x}\left[T_{(0, x+\xi)}>t\right] \\
& \geqq \int_{F(x+\xi)} d y \int_{[0, \infty)} d \xi P_{y}\left[T_{F(x+\xi)}>t\right] \frac{\partial}{\partial \xi} P_{x}\left[T_{(0, x+\xi)}>t\right] \\
& =-\int_{0}^{\infty} d \xi P_{x}\left[T_{(0, x+\xi)}>t\right] \frac{\partial}{\partial \xi} Q_{F(x+\xi)}(t) \tag{5.32}
\end{align*}
$$

Since $F$ is horn-shaped $-\frac{\partial}{\partial \xi} Q_{F(x+\xi)}(t)$ is positive. Hence by Lemma 5.9 and further integrations by parts

$$
\begin{align*}
\int_{F(x)} d y P_{(x, y)}\left[T_{F}>t\right] \geqq & \int_{0}^{\infty} d \xi\left\{1-\int_{\xi}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime}-\int_{x}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime}\right\} \cdot-\frac{\partial}{\partial \xi} Q_{F(x+\xi)}(t) \\
= & Q_{F(x)}(t)+\int_{[0, \infty)} d \xi\left\{Q_{F(x+\xi)}(t)-Q_{F(x)}(t)\right\} \rho(\xi ; t) \\
& -Q_{F(x)}(t) \int_{x}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime} \tag{5.33}
\end{align*}
$$

Hence

$$
\begin{align*}
\int_{0}^{\infty} d x \int_{F(x)} d y P_{(x, y)}\left[T_{F}>t\right] \geqq & \int_{0}^{\infty} d x Q_{F(x)}(t)-\int_{0}^{\infty} d \xi \rho(\xi ; t) \int_{0}^{\xi} d x Q_{F(x)}(t) \\
& -\int_{0}^{\infty} d x Q_{F(x)}(t) \int_{x}^{\infty} \rho\left(\xi^{\prime} ; t\right) d \xi^{\prime} \tag{5.34}
\end{align*}
$$

A further integration by parts completes the proof.
It is possible to obtain the precise asymptotic behaviour of $Z(t), Q(t)$ as $t \downarrow 0$ for other unbounded regions. In particular we have the following.
Theorem 5.10. Let $S_{\lambda} \subset \mathbb{R}^{2}$ be the complement of the range of the curve $\gamma:[0, \infty)$ $\rightarrow \mathbb{R}^{2}$ defined by

$$
\begin{equation*}
\gamma(s)=\left(s^{2} \cos s, s^{2} \sin s\right), \quad 0<\lambda<1 \tag{5.35}
\end{equation*}
$$

Then $S_{\lambda}$ is open and dense in $\mathbb{R}^{2}$ and has infinite volume. Furthermore

$$
\begin{align*}
\lim _{t \downarrow 0} Q_{S_{\lambda}}(t) t^{\lambda /(1-\lambda)}= & \frac{2}{\pi} \lambda^{2 \lambda /(1-\lambda)}\left(4^{1 /(1-\lambda)}-1\right) \Gamma(1 /(1-\lambda)) \zeta(2 /(1-\lambda)),  \tag{5.36}\\
\lim _{t \downarrow 0} Z_{S_{\lambda}}(t) t^{1 /(1-\lambda)}= & \frac{1}{\pi^{1 / 2}(1-\lambda)} \lambda^{(1+\lambda) /(1-\lambda)} 2^{(3 \lambda-1) /(1-\lambda)} \Gamma((1+\lambda) /(2-2 \lambda)) \\
& \cdot \zeta((1+\lambda) /(1-\lambda)), \tag{5.37}
\end{align*}
$$

where $\zeta, \Gamma$ are the Riemann function, gamma function respectively.
Sketch of Proof. By "unrolling" the spiral one obtains a (one-sided) horn-shaped region of the form (5.7) with $f(x) \sim(s+2 \pi)^{2}-s^{\lambda} \sim 2 \pi \lambda s^{\lambda-1}$ as $s \uparrow \infty$ and $d x$ $\sim\left(s^{2 \lambda}+\lambda^{2} s^{2 \lambda-2}\right)^{1 / 2} d s \sim s^{\lambda} d s$ as $s \uparrow \infty$. Formulas (5.8) and (5.16) yield (5.36) and (5.37) respectively.

Theorem 5.11. Let $F$ be horn-shaped in $\mathbb{R}^{2}$. Then for $0 \leqq t<\infty$ the following are equivalent. (i) $Q_{F}(s)<\infty$ for all $s>t$, (ii) $Z_{F}(s)<\infty$ for all $s>t$.
Proof. Because of Lemma 2.6 we only have to prove that $Q_{F}(s)<\infty$ for $s>t$ implies $Z_{F}(s)<\infty$ for $s>t$. From (5.8) and [2] we obtain (for one-sided hornshaped regions in $\mathbb{R}^{2}$ )

$$
\begin{align*}
Z_{F}(s) & \leqq \frac{1}{(4 \pi s)^{1 / 2}} \int_{0}^{\infty} d x \sum_{k=1}^{\infty} e^{-s \pi^{2} k^{2} / f^{2}(x)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{1}{(4 \pi s)^{1 / 2}} \int_{t^{1 / 2}}^{\infty} d x \sum_{k=1}^{\infty} e^{-s \pi^{2} k^{2} / f^{2}(x)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{2}{(4 \pi s)^{1 / 2}} \int_{t^{1 / 2}}^{\infty} d x \sum_{k=1,3, \ldots} e^{-s \pi^{2} k^{2} / f^{2}(x)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{2}{(4 \pi s)^{1 / 2}} \int_{t^{1 / 2}}^{\infty} d x \sum_{k=1,3} e^{-(s+t) \pi^{2} k^{2} /\left(2 f^{2}(x)\right)-(s-t) \pi^{2} k^{2} /\left(2 f^{2}(x)\right)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{4}{(4 \pi s)^{1 / 2}} \frac{1}{(s-t) \pi^{2}} \int_{t^{1 / 2}}^{\infty} d x f^{2}(x) \sum_{k=1,3, \ldots} k^{-2} e^{-(s+t) \pi^{2} k^{2} /\left(2 f^{2}(x)\right)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{f\left(t^{1 / 2}\right)}{(16 \pi s)^{1 / 2}(s-t)} \int_{0}^{\infty} d x f(x) \sum_{k=1,3} \frac{8}{\pi^{2} k^{2}} e^{-(s+t) \pi^{2} k^{2} /\left(2 f^{2}(x)\right)} \\
& \leqq \frac{1}{4 \pi s} \int_{0}^{t^{1 / 2}} d x f(x)+\frac{f\left(t^{1 / 2}\right)}{(16 \pi s)^{1 / 2}(s-t)}\left\{Q_{F}((s+t) / 2)+4 \int_{0}^{t^{1 / 2}} d x f(x)\right\} . \tag{5.38}
\end{align*}
$$

This is finite since $f$ is integrable at 0 and $Q_{F}((s+t) / 2)$ is finite by $(s+t) / 2>t$.

## 6. Bounded Regions with $R$-smooth Boundaries

Before we state the main theorem of this section we make the following definition.
Definition 6.1. A boundary $\partial D$ of an open set $D$ in $\mathbb{R}^{m}, m=2,3, \ldots$ is $R$-smooth if for each point $x_{0} \in \partial D$ there exist two open balls $B_{1}, B_{2}$ with radius $R$ such that $B_{1} \subset D, B_{2} \subset \mathbb{R}^{m} \backslash(D \cup \partial D), \partial B_{1} \cap \partial B_{2}=x_{0}$.

Theorem 6.2. Let $D$ be an open, bounded and connected set in $\mathbb{R}^{m}, m=2,3, \ldots$ with $R$-smooth boundary $\partial D$. Then for all $t>0$

$$
\begin{gather*}
\left|Q_{D}(t)-|D|+\frac{2}{\pi^{1 / 2}} t^{1 / 2}\right| \partial D\left|\left|\leqq 10^{m}\right| D\right| t R^{-2}  \tag{6.1a}\\
\left|(4 \pi t)^{m / 2} Z_{D}(t)-|D|+\frac{\pi^{1 / 2}}{2} t^{1 / 2}\right| \partial D\left|\left|\leqq 2^{m} m^{4}\right| D\right| t R^{-2} \tag{6.1b}
\end{gather*}
$$

where $|\partial D|$ is the area of the boundary $\partial D$.
The proof of (6.1b) has been given in [2]. Here we will prove (6.1a). In Lemma 6.3 and Corollary 6.4 we will use Levy's maximal inequality for brownian motion to obtain estimates for $P_{x}\left[T_{B}>t\right]$ where $B$ is an open ball in $\mathbb{R}^{m}$. In Lemmas 6.5 and 6.6 we obtain a lower bound and upper bound respectively for $x$ near the boundary $\partial D$. In Lemma 6.7 we recall a result on areas of parallel surfaces. Then we complete the proof of (6.1a).

Lemma 6.3. Let $B$ be an open ball in $\mathbb{R}^{m}$ with radius $R$ and centre 0 . Then

$$
\begin{equation*}
1 \geqq P_{0}\left[T_{B}>t\right] \geqq 1-\frac{2}{\Gamma(m / 2)} \int_{R^{2} /(4 t)}^{\infty} y^{(m-2) / 2} e^{-y} d y \tag{6.2}
\end{equation*}
$$

Proof. By Levy's maximal inequality (Theorem 3.6.5 of [19])

$$
\begin{align*}
P_{0}\left[T_{B}<t\right] & =P_{0}\left[\max _{0 \leqq \tau \leqq t}|x(\tau)|>R\right] \leqq 2 P_{0}[|x(t)|>R] \\
& =\frac{2}{(4 \pi t)^{m / 2}} \int_{|x|>R} d x e^{-|x|^{2} /(4 t)} . \tag{6.3}
\end{align*}
$$

Corollary 6.4. Let B be as in Lemma 6.3. Then

$$
\begin{equation*}
P_{0}\left[T_{B}<t\right] \leqq 2^{1+m / 2} e^{-R^{2} /(8 t)} \tag{6.4}
\end{equation*}
$$

Lemma 6.5. Let $D$ be open in $\mathbb{R}^{m}$ with $R$-smooth boundary $\partial D$. Let $x \in D$ such that $d(x)<R$. Then

$$
\begin{align*}
P_{x}\left[T_{D}>t\right] \geqq & \int_{0}^{d(x)} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)}\left\{1-\frac{2}{\Gamma((m-1) / 2)} \int_{R(d(x)-\xi) /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2}\right\} \\
& -\frac{2}{\Gamma(m / 2)} \int_{R d(x) /(4 t)}^{\infty} d y e^{-y} y^{(m-2) / 2} . \tag{6.5}
\end{align*}
$$

Proof. Since $\partial D$ is $R$-smooth and $d(x)<R$ there exists a ball $B$ with radius $R$ and centre 0 such that $0=(-(R-d(x)), 0)$ and $x=(0,0)$ in a cartesian frame
$\left(x_{1}, y\right),\left(x_{1} \in \mathbb{R}, y \in \mathbb{R}^{m-1}\right)$. Let $\partial_{1} B=\left\{z \in \partial B \mid x_{1}(z) \geqq 0\right\}, \partial_{2} B=\partial B \backslash \partial_{1} B$. Then by monotonicity

$$
\begin{align*}
P_{x}\left[T_{D}>t\right] \geqq & P_{x}\left[T_{B}>t\right] \geqq P_{(0,0)}\left[x(\cdot) \text { does not hit } \partial_{1} B \text { up to time } t\right] \\
& -P_{(0,0)}\left[x(\cdot) \text { does hit } \partial_{2} B \text { in }[0, t]\right] . \tag{6.6}
\end{align*}
$$

Since the distance from $(0,0)$ to $\partial_{2} B$ is bounded from below by $(\operatorname{Rd}(x))^{1 / 2}$ we have by Lemma 6.3

$$
\begin{align*}
P_{(0,0)}\left[x(\cdot) \text { does hit } \partial_{2} B \text { in }[0, t]\right] & \leqq P_{(0,0)}\left[\max _{0 \leqq \tau \leqq t}|x(\tau)|^{2}>R d(x)\right] \\
& \leqq \frac{2}{\Gamma(m / 2)} \int_{R d(x) /(4 t)}^{\infty} d y e^{-y} y^{(m-2) / 2} \tag{6.7}
\end{align*}
$$

Let $x(\cdot)=\left(x_{1}(\cdot), y(\cdot)\right)$ where $x_{1}(\cdot)$ is a brownian motion along the $x_{1}$ axis with $x_{1}(0)=0$ and $y(\cdot)$ is an independent brownian motion in $\mathbb{R}^{m-1}$ with $y(0)=0$. The probability density $\rho(\xi ; t)$ of the random variable $\max _{0 \leqq \tau \leqq t} x_{1}(\tau)$ is given by (5.24). Hence by mononicity and Lemma 6.3

$$
\begin{align*}
& P_{(0,0)}\left[x(\cdot) \text { does not hit } \partial_{1} B \text { up to time } t\right] \\
& \quad \geqq \int_{0}^{d(x)} d \xi \rho(\xi ; t) P_{0}\left[\max _{0 \leqq \tau \leqq t}|y(\tau)|^{2}<R^{2}-(R-d(x)+\xi)^{2}\right] \\
& \quad \geqq \int_{0}^{d(x)} d \xi \rho(\xi ; t) P_{0}\left[\max _{0 \leqq \tau \leqq t}|y(\tau)|^{2}<R(d(x)-\xi)\right] \\
& \quad \geqq \int_{0}^{d(x)} d \xi \rho(\xi ; t)\left\{1-\frac{2}{\Gamma((m-1) / 2)} \int_{R(d(x)-\xi) /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2}\right\} . \tag{6.8}
\end{align*}
$$

Lemma 6.6. Let $D$ be open in $\mathbb{R}^{m}$ with $R$-smooth boundary $\partial D$. Let $x \in D$ such that $d(x)<R$. Then

$$
\begin{equation*}
P_{x}\left[T_{D}>t\right] \leqq \int_{0}^{d(x)} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)}+\frac{4(m-1) t^{1 / 2}}{\pi^{1 / 2} R} e^{-d^{2}(x) /(4 t)}+2^{9 / 2} t /\left(e R^{2}\right) \tag{6.9}
\end{equation*}
$$

Proof. Since $\partial D$ is $R$-smooth there exists a closed ball $B$ with radius $R$ in the complement of $D$ such that the distance from $x$ to $\partial B$ is equal to $d(x)$. Define a cartesian frame $\left(x_{1}, y\right),\left(x_{1} \in \mathbb{R}, y \in \mathbb{R}^{m-1}\right)$ such that $x=(0,0)$ and the centre of $B$ is given by $(d(x)+R, 0)$. Then by monotonicity and the decomposition of the brownian motion (as in the proof of Lemma 6.5)

$$
\begin{align*}
P_{x}\left[T_{D}>t\right] \leqq & P_{x}\left[T_{\mathbb{R} m \backslash B}>t\right] \leqq \int_{0}^{d(x)} \rho(\xi ; t) d \xi+\int_{d(x)+R}^{\infty} \rho(\xi ; t) d \xi \\
& +\int_{d(x)}^{d(x)+R} \rho(\xi ; t) d \xi P_{0}\left[\max _{0 \leqq \tau \leqq t}|y(\tau)|^{2}>R^{2}-(R-\xi+d(x))^{2}\right] . \tag{6.10}
\end{align*}
$$

## By Lemma 6.3

$$
\begin{align*}
& \int_{d(x)}^{d(x)+R} \rho(\xi ; t) d \xi P_{0}\left[\max _{0 \leqq \tau \leqq t}|y(\tau)|^{2}>R^{2}-(R-\xi+d(x))^{2}\right] \\
& \quad \leqq \int_{d(x)}^{d(x)+R} \rho(\xi ; t) d \xi P_{0}\left[\max _{0 \leq \tau \leq t}|y(\tau)|^{2}>R(\xi-d(x))\right] \\
& \quad \leqq \int_{d(x)}^{d(x)+R} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-d^{2}(x) /(4 t)} \frac{2}{\Gamma((m-1) / 2)} \int_{R(\xi-d(x)) /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2} \\
& \quad \leqq \frac{2}{(\pi t)^{1 / 2}} \frac{e^{-d^{2}(x) /(4 t)}}{\Gamma((m-1) / 2)} \int_{0}^{\infty} d \xi \int_{R \xi /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2} . \tag{6.11}
\end{align*}
$$

Furthermore by Corollary 6.4 for $m=1$

$$
\begin{equation*}
\int_{d(x)+R}^{\infty} \rho(\xi ; t) d \xi \leqq 2^{3 / 2} e^{-R^{2} /(8 t)} \leqq 2^{9 / 2} t /\left(e R^{2}\right) \tag{6.12}
\end{equation*}
$$

Lemma 6.7. Let $D$ be an open, bounded and connected set in $\mathbb{R}^{m}, m=2,3, \ldots$ with $R$-smooth boundary $\partial D$. Let $\partial D_{q}$ denote the boundary of the set $\{x \in D \mid d(x)$ $>q\}$ and let $\left|\partial D_{q}\right|$ denote its area. Then

$$
\begin{equation*}
|\partial D|\left(\frac{R-q}{R}\right)^{m-1} \leqq\left|\partial D_{q}\right| \leqq|\partial D|\left(\frac{R}{R-q}\right)^{m-1}, \quad 0 \leqq q<R \tag{6.13}
\end{equation*}
$$

Proof. See Lemma 5 of [2].
Proof of Theorem 6.2. By Lemmas 6.6 and 6.7 we have

$$
\begin{align*}
Q_{D}(t)= & |D|-\int_{D} d x P_{x}\left[T_{D}<t\right] \\
\leqq & |D|-\int_{\{x \in D: d(x)<R / 2\}} d x P_{x}\left[T_{D}<t\right] \\
\leqq & |D|-\int_{0}^{R / 2}\left|\partial D_{q}\right| d q\left\{\int_{q}^{\infty} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)}-\frac{4(m-1) t^{1 / 2}}{\pi^{1 / 2} R} e^{-q^{2} /(4 t)}-\frac{2^{9 / 2} t}{e R^{2}}\right\} \\
\leqq & |D|-\int_{0}^{R / 2}|\partial D|\left(1-\frac{q}{R}\right)^{m-1} d q \int_{q}^{\infty} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)} \\
& +2^{m-1}|\partial D| \int_{0}^{\infty} d q \frac{4(m-1) t^{1 / 2}}{\pi^{1 / 2} R} e^{-q^{2} /(4 t)}+|D| 2^{9 / 2} t /\left(e R^{2}\right) \\
\leqq & |D|-\int_{0}^{\infty} d q|\partial D|(1-2(m-1) q / R) \int_{q}^{\infty} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)} \\
& +2^{m+1}(m-1)|\partial D| t / R+2^{9 / 2}|D| t /\left(e R^{2}\right) \\
= & |D|-\frac{2 t^{1 / 2}}{\pi^{1 / 2}}|\partial D|+\frac{|\partial D| t}{R} 2(m-1)\left(1+2^{m}\right)+2^{9 / 2}|D| t /\left(e R^{2}\right) \tag{6.14}
\end{align*}
$$

From Lemma 6.7 we obtain by integrating with respect to $q$

$$
\begin{equation*}
|\partial D| \leqq \frac{m|D|}{R} \tag{6.15}
\end{equation*}
$$

The upper bound in Theorem 6.2 follows from (6.14) and (6.15).
To prove the lower bound in Theorem 6.2 we use Corollary 6.4 and Lemmas 6.5 and 6.7 . We obtain

$$
\begin{align*}
Q_{D}(t) \geqq & \int_{\{x \in D: d(x)>R / 2\}} d x\left(1-2^{1+m / 2} e^{-R^{2} /(32 t)}\right)+\int_{\{x \in D: d(x) \leqq R / 2\}} d x P_{x}\left[T_{D}>t\right] \\
\geqq & |D|-|D| t 2^{6+m / 2} /\left(e R^{2}\right)-\int_{\{x \in D: d(x) \leqq R / 2\}} d x \int_{d(x)}^{\infty} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)} \\
& -\frac{2}{\Gamma((m-1) / 2)} \int_{\{x \in D: d(x) \leqq R / 2\}} \int_{0}^{d(x)} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)} \int_{R(d(x)-\xi) /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2} \\
& -\frac{2}{\Gamma(m / 2)} \int_{\{x \in D: d(x) \leqq R / 2\}} \int_{R d(x) /(4 t)}^{\infty} d y e^{-y} y^{(m-2) / 2} \\
\geqq & |D|-|D| t 2^{6+m / 2} /\left(e R^{2}\right)-\int_{0}^{\infty} d q|\partial D|\left(1+(m-1) 2^{m-1} q / R\right) \int_{q}^{\infty} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2 / /(4 t)}} \\
& -\frac{2}{\Gamma((m-1) / 2)} \cdot 2^{m-1}|\partial D| \int_{0}^{\infty} d q \int_{0}^{q} \frac{d \xi}{(\pi t)^{1 / 2}} e^{-\xi^{2} /(4 t)} \int_{R(q-\xi) /(4 t)}^{\infty} d y e^{-y} y^{(m-3) / 2} \\
& -\frac{2}{\Gamma(m / 2)} \cdot 2^{m-1}|\partial D| \int_{0}^{\infty} d q \int_{R q /(4 t)}^{\infty} d y e^{-y} y^{(m-2) / 2} \\
= & |D|-\frac{2 t^{1 / 2}}{\pi^{1 / 2}}|\partial D|-|D| t 2^{6+m / 2} /\left(e R^{2}\right)-\frac{|\partial D| t}{R}(9 m-5) 2^{m-1} . \tag{6.16}
\end{align*}
$$

The lower bound in Theorem 6.2 follows from (6.15) and (6.16).
Corollary 6.8. Let $D$ be an open, bounded and connected set in $\mathbb{R}^{m}, m=2,3, \ldots$ with $R$-smooth boundary $\partial D$. Let $h:[0, \infty) \rightarrow R$ be $C^{1}$ and $h(0)=0$. Let $q: \bar{D}$ $\times[0, \infty) \rightarrow R$ be the (classical) solution of

$$
\begin{align*}
\Delta q & =\frac{\partial q}{\partial t} \quad \text { on } D \times(0, \infty)  \tag{6.17}\\
q(x ; t) & =h(t), \quad x \in \partial D, \quad t \geqq 0  \tag{6.18}\\
q(x ; 0) & =0, \quad x \in D \tag{6.19}
\end{align*}
$$

where $\Delta$ is the laplacian. Then for $t>0$

$$
\begin{equation*}
\left|\int_{D} q(x ; t) d x-\frac{|\partial D|}{\pi^{1 / 2}} \int_{0}^{t} h(\tau)(t-\tau)^{-1 / 2} d \tau\right| \leqq 10^{m} \frac{|D|}{R^{2}} \int_{0}^{t}\left|h^{\prime}(\tau)\right|(t-\tau) d \tau . \tag{6.20}
\end{equation*}
$$

Proof. The solution of (6.17)-(6.19) is given by

$$
\begin{equation*}
q(x ; t)=\int_{0}^{t} h^{\prime}(\tau) P_{x}\left[T_{D}<t-\tau\right] d \tau . \tag{6.21}
\end{equation*}
$$

Hence by Fubini's theorem

$$
\begin{equation*}
\int_{D} q(x ; t) d x=\int_{0}^{t} h^{\prime}(\tau)\left(|D|-Q_{D}(t-\tau)\right) d \tau . \tag{6.22}
\end{equation*}
$$

Corollary 6.8 follows from (6.22) and Theorem 6.2.

## References

1. Berg, M. van den: On the spectrum of the Dirichlet Laplacian for horn-shaped regions in $\mathbb{R}^{n}$ with infinite volume. J. Funct. Anal. 58, 150-156 (1984)
2. Berg, M. van den: On the asymptotics of the heat equation and bounds on traces associated with the Dirichlet Laplacian. J. Funct. Anal. 71, 279-293 (1987)
3. Berg, M. van den, Srisatkunarajah, S.: Heat flow and brownian motion for a region in $\mathbb{R}^{2}$ with a polygonal boundary. (Preprint 1989)
4. Berg, M. van den: Gaussian bounds for the Dirichlet heat kernel. J. Funct. Anal. (to appear)
5. Carslaw, H.S., Jaeger, J.C.: Conduction of heat in solids. Oxford: Clarendon Press 1986
6. Davies, E.B.: Some norm bounds and quadratic form inequalities for Schrödinger operators, II. J. Oper. Theory 12, 177-196 (1984)
7. Davies, E.B.: Trace properties of the Dirichlet Laplacian. Math. Zeit. 188, 245-251 (1985)
8. Davies, E.B.: One-parameter semigroups. New York: Academic Press 1980
9. Davies, E.B.: Heat kernels and spectral theory. Cambridge: Cambridge University Press 1989
10. Davies, E.B., Mandouvalos, N.: Heat kernel bounds on manifolds with cusps. J. Funct. Anal. 75, 311-322 (1987)
11. Davies, E.B., Simon, B.: Ultracontractivity and the heat kernel for Schrödinger operators and Dirichlet Laplacians. J. Funct. Anal. 59, 335-395 (1984)
12. Edmunds, D.E., Evans, W.D.: Spectral theory and differential operators. Oxford: Clarendon Press 1987
13. Fleckinger, J.: Asymptotic distribution of eigenvalues of elliptic operators on unbounded domains. In: Everitt, W.N., Sleeman, B.D. (eds.) Ordinary and partial differential equations. Proceedings, Dundee 1980. (Lect. Notes Math., vol. 846, pp. 119-128) Berlin Heidelberg New York: Springer 1981
14. Gradshteyn, I.S., Ryzhik, I.M.: Table of integrals, series and products, New York: Academic Press 1980
15. Maz'ja, V.G.: On ( $p, l$-capacity, imbedding theorems and the spectrum of a selfadjoint elliptic operator. Math. USSR, Izv. 7, 357-387 (1973)
16. Rosen, J.: Sobolev inequalities for weight spaces and supercontractivity. Trans. Am. Math. Soc. 222, 367-376 (1976)
17. Rozenbljum, V.G.: The eigenvalues of the first boundary value problem in unbounded domains. Math. USSR, Sb. 18, 235-248 (1972)
18. Rozenbljum, V.G.: The calculation of the spectral properties for the Laplace operator in domains of infinite measure. In: Smirnov, V.I. (ed.) Problems of mathematical analysis. No. 4: Integral and differential operators. Differential equations. Leningrad: Izdat. Leningrad Univ. 1973
19. Simon, B.: Functional integration and quantum physics. New York: Academic Press 1979
20. Simon, B.: Classical boundary conditions as a technical tool in modern mathematical physics. Adv. Math. 30, 377-385 (1978)
21. Simon, B.: Some quantum operators with discrete spectrum but classically continuous spectrum. Ann. Phys. 146, 209-220 (1983)
22. Simon, B.: Non-classical eigenvalue asymptotics. J. Funct. Anal. 53, $84-98$ (1983)
23. Tamura, H.: The asymptotic distribution of eigenvalues of the Laplace operator in an unbounded domain. Nagoya Math. J. 60, 7-33 (1976)

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