# How to initialise a second class particle? Joint with Attila László Nagy 

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Advances in Last Passage Percolation 26 June, 2019.

The models

Simple exclusion<br>Zero range<br>Bricklayers

Hydrodynamics

The second class particle

Ferrari-Kipnis for TASEP

Let's generalise

## Totally asymmetric simple exclusion



Bernoulli( $\varrho$ ) distribution; $\omega_{i}=0$ or 1.

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The Bernoulli( $\varrho$ ) distribution is time-stationary for any $(0 \leq \varrho \leq 1)$. Any translation-invariant stationary distribution is a mixture of Bernoullis.

## Totally asymmetric zero range process

$$
\omega_{i} \in \mathbb{Z}^{+}
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Particles jump to the right with rate $r\left(\omega_{i}\right) \quad$ ( $r$ non-decreasing).

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## Totally asymmetric zero range process

Extremal translation-invariant stationary distributions are still product, and rather explicit in terms of $r(\cdot)$.

Two special cases:

- $r\left(\omega_{i}\right)=\mathbf{1}\left\{\omega_{i}>0\right\}$ : classical zero range; $\omega_{i} \sim \operatorname{Geom}(\theta)$.
- $r\left(\omega_{i}\right)=\omega_{i}$ : independent walkers; $\omega_{i} \sim \operatorname{Poi}(\theta)$.


## Totally asymmetric bricklayers process



## Totally asymmetric bricklayers process


a brick is added with rate $\left[r\left(\omega_{i}\right)+r\left(-\omega_{i+1}\right)\right]$

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(r(\omega) \cdot r(1-\omega)=1 ; \quad r \text { non-decreasing }) .
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Extremal translation-invariant stationary distributions are still product, and rather explicit in terms of $r(\cdot)$.

A special case: $r\left(\omega_{i}\right)=\mathrm{e}^{\beta \omega_{i}}: \omega_{i} \sim \operatorname{discrete} \operatorname{Gaussian}\left(\frac{\theta}{\beta}, \frac{1}{\sqrt{\beta}}\right)$.

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## Rescaled version: rarefaction fan



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## The second class particle

States $\omega$ and $\eta$ only differ at one site.


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## Growth on the right: rate $\leq$ rate

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## Growth on the right: rate $\leq$ rate with rate:

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Growth on the left: rate $\geq$ rate


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Growth on the left: rate $\geq$ rate with rate-rate:


A single discrepancy $\uparrow$, the second class particle, is conserved. Its position at time $t$ is $Q(t)$.

## Ferrari-Kipnis '95 for TASEP

Blue TASEP $\omega$ :
Bernoulli( $\varrho$ ) for sites $\{\ldots,-2,-1,0\}$,
Bernoulli $(\lambda)$ for sites $\{1,2,3, \ldots\}$.
Black TASEP $\eta$ :
Bernoulli $(\varrho)$ for sites $\{\ldots,-3,-2,-1\}$,
Bernoulli $(\lambda)$ for sites $\{0,1,2, \ldots\}$.

$h_{i}(t), g_{i}(t)$ are the respective numbers of particles jumping over the edge $(i, i+1)$ by time $t(i>0)$.

## Ferrari-Kipnis '95 for TASEP, Part 1

First realization:


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- $\omega_{i}(0)=\eta_{i}(0) \sim \operatorname{Bernoulli}(\varrho)$ for $i<0$



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- $\omega_{i}(0)=\eta_{i}(0) \sim \operatorname{Bernoulli}(\varrho)$ for $i<0$
- $\left(\omega_{0}(0), \eta_{0}(0)\right)=(0,0)$ w. prob. $1-\varrho$ $\left(\omega_{0}(0), \eta_{0}(0)\right)=(1,0)$ w. prob. $\varrho-\lambda$ $\left(\omega_{0}(0), \eta_{0}(0)\right)=(1,1)$ w. prob. $\lambda$



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- $\omega_{i}(0)=\eta_{i}(0) \sim \operatorname{Bernoulli}(\lambda)$ for $i>0$



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- $\omega_{i}(0)=\eta_{i}(0) \sim \operatorname{Bernoulli}(\lambda)$ for $i>0$

$\mathbf{E} h_{i}(t)-\mathbf{E} g_{i}(t)=\mathbf{E}\left(h_{i}(t)-g_{i}(t)\right)=(\varrho-\lambda) \cdot \mathbf{P}\{Q(t)>i\}$.


## Ferrari-Kipnis '95 for TASEP, Part 2

Second realization:

$$
\omega_{i}(t) \equiv \eta_{i-1}(t) \quad \forall i, \forall t .
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$\mathbf{E} h_{i}(t)-\mathbf{E} g_{i}(t)=\mathbf{E}\left(h_{i}(t)-g_{i}(t)\right)=\mathbf{E}\left(\eta_{i}(t)-\eta_{i}(0)\right)=\mathbf{E} \eta_{i}(t)-\mathbf{E} \eta_{i}(0)$.

## Ferrari-Kipnis '95 for TASEP

Thus,

$$
\begin{aligned}
& \mathbf{E} h_{i}(t)-\mathbf{E} g_{i}(t)=\mathbf{E}\left(h_{i}(t)-g_{i}(t)\right)=(\varrho-\lambda) \cdot \mathbf{P}\{Q(t)>i\}, \\
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$$
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$$

Combine with hydrodynamics to conclude

$$
\frac{Q(t)}{t} \Rightarrow \begin{cases}\text { shock velocity } & \text { in a shock } \\ \mathrm{U}\left(H^{\prime}(\varrho), H^{\prime}(\lambda)\right) & \text { in a rarefaction wave. }\end{cases}
$$

## Let's generalise

Other models have more than 0 or 1 particles per site. How do we start the second class particle?

- Shall we do

$$
\left.()^{0}\right) \quad \operatorname{Or} ? \quad \text { Or } e_{e}^{\bullet} \text { ? Or } \ldots
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## Let's generalise

Other models have more than 0 or 1 particles per site. How do we start the second class particle?

- Shall we do
- Recall for TASEP we increased $\lambda$ to $\varrho$ by adding or not adding a $2^{\text {nd }}$ class particle.
$\left(\omega_{0}(0), \eta_{0}(0)\right)=(0,0)$ w. prob. $1-\varrho$
$\left(\omega_{0}(0), \eta_{0}(0)\right)=(1,0)$ w. prob. $\varrho-\lambda$
$\left(\omega_{0}(0), \eta_{0}(0)\right)=(1,1)$ w. prob. $\lambda$
$\varrho_{i}$



## Let's generalise: problems with coupling

Fix $\lambda<\varrho \leq \lambda+1$. Is there a joint distribution of $\left(\omega_{0}, \eta_{0}\right)$ such that

- the first marginal is $\omega_{0} \sim$ stati. $\mu^{\varrho}$;
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Proposition

- Of course for Bernoulli (TASEP).


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- No for Poisson (indep. walkers with $r\left(\omega_{i}\right)=\omega_{i}$ ).
- Yes for discrete Gaussian (bricklayers with $r\left(\omega_{i}\right)=\mathrm{e}^{\beta \omega_{i}}$ ).


## Let's generalise

Keep calm and couple anyway.
Find a coupling measure $\nu$ with

- first marginal $\omega_{0} \sim$ stati. $\mu^{\varrho}$;
- second marginal $\eta_{0} \sim$ stati. $\mu^{\lambda}$;
- zero weight whenever $\omega_{0} \notin\left\{\eta_{0}, \eta_{0}+1\right\}$.

Not many choices:

$$
\begin{aligned}
\nu(x, x) & =\mu^{\varrho}\{-\infty \ldots x\}-\mu^{\lambda}\{-\infty \ldots x-1\} \\
\nu(x+1, x) & =\mu^{\lambda}\{-\infty \ldots x\}-\mu^{\varrho}\{-\infty \ldots x\} \\
\nu & =\text { zero elsewhere }
\end{aligned}
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- Bad news: $\nu(x, x)$ can be negative (e.g., Geom., Poi).


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- Good news: Who cares? No $2^{\text {nd }}$ class particle there.


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- Good news: $\nu(x+1, x) \geq 0$ (attractivity).


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- Bad news: $\nu(x, x)$ can be negative (e.g., Geom., Poi).
- Good news: Who cares? No $2^{\text {nd }}$ class particle there.
- Good news: $\nu(x+1, x) \geq 0$ (attractivity).

We can still use the signed measure $\nu$ formally, as we only care about $\nu(x+1, x)$. Scale this up to get the initial distribution at the site of the second class particle:

$$
\mu\left(\omega_{0}, \eta_{0}\right)=\mu\left(\eta_{0}+1, \eta_{0}\right)=\frac{\nu\left(\eta_{0}+1, \eta_{0}\right)}{\sum_{x} \nu(x+1, x)}=\frac{\nu\left(\eta_{0}+1, \eta_{0}\right)}{\varrho-\lambda} .
$$

## Let's generalise

$$
\mu\left(\omega_{0}, \eta_{0}\right)=\frac{\nu\left(\eta_{0}+1, \eta_{0}\right)}{\varrho-\lambda}
$$

- is a proper probability distribution;
- actually agrees with the coupling measure $\nu$ conditioned on a $2^{\text {nd }}$ class particle when $\nu$ behaves nicely (Bernoulli, discr.Gaussian);


## Let's generalise

Theorem
Starting in

$$
\begin{gathered}
\bigotimes_{i<0} \mu_{i}^{\varrho} \otimes \mu_{0} \otimes \bigotimes_{i>0} \mu_{i}^{\lambda}, \\
\lim _{N \rightarrow \infty} \mathbf{P}\left\{\frac{Q(N T)}{N}>X\right\}=\frac{\varrho(X, T)-\lambda}{\varrho-\lambda}
\end{gathered}
$$

where $\varrho(X, T)$ is the entropy solution of the hydrodynamic equation with initial data
$\varrho$ on the left
$\lambda$ on the right.

## What do we have?

$$
\lim _{N \rightarrow \infty} \mathbf{P}\left\{\frac{Q(N T)}{N}>X\right\}=\frac{\varrho(X, T)-\lambda}{\varrho-\lambda}
$$

$\rightsquigarrow$ The solution $\varrho(X, T)$ is the distribution of the velocity for $Q$.

- Shock: distribution is step function, velocity is deterministic (LLN).
- Rarefaction wave: distribution is continuous, velocity is random (e.g., Uniform for TASEP).


## A fun model (B., A.L. Nagy, I. Tóth, B. Tóth)

$$
\begin{array}{rlrl}
\omega_{i}=-1,0,1 ; & & \\
(0,-1) & \rightarrow(-1,0) & & \text { with rate } \frac{1}{2} \\
(1,0) & \rightarrow(0,1) & & \text { with rate } \frac{1}{2} \\
(1,-1) & \rightarrow(0,0) & & \text { with rate } 1 \\
(0,0) & \rightarrow(-1,1) & & \text { with rate } c .
\end{array}
$$

## A fun model (B., A.L. Nagy, I. Tóth, B. Tóth)

Hydrodynamic flux $H(\varrho)$, for certain $c$ :


## A fun model (B., A.L. Nagy, I. Tóth, B. Tóth)

Here is what can happen (R: rarefaction wave, $\mathbf{S}$ : Shock):


## A fun model (B., A.L. Nagy, I. Tóth, B. Tóth)

Examples for $\varrho(T, X)$ :


$$
\lim _{N \rightarrow \infty} \mathbf{P}\left\{\frac{Q(N T)}{N}>X\right\}=\frac{\varrho(X, T)-\lambda}{\varrho-\lambda}
$$

$\rightsquigarrow$ The solution $\varrho(X, T)$ is the distribution of the velocity for $Q$.
I haven't seen a walk with a random velocity of mixed distribution before.

## Storytelling...

$$
\mu\left(\omega_{0}, \eta_{0}\right)=\frac{\nu\left(\eta_{0}+1, \eta_{0}\right)}{\varrho-\lambda}
$$

In the $1 / 3$-fluctuations papers (B., J. Komjáthy, T. Seppäläinen)
we had to start the second class particle in a $\varrho=\lambda$ flat environment. We came up with a measure $\hat{\mu}$ for this which worked nicely with our formulas. But at that time we had no idea why.

As it turns out: $\hat{\mu}=\lim _{\lambda \nearrow \varrho} \mu$.

## Symmetric case

Everything works with partially asymmetric models (allow left jumps too).

In fact everything works for symmetric models as well. The hydrodynamic scaling is diffusive there with the limit being of heat equation type. In this case:

## Symmetric case

Theorem (Symmetric version)
Starting in

$$
\begin{gathered}
\bigotimes \mu_{i}^{\varrho} \otimes \mu_{0} \otimes \bigotimes_{i>0} \mu_{i}^{\lambda}, \\
\lim _{N \rightarrow \infty} \mathbf{P}\left\{\frac{Q(N T)}{\sqrt{N}}>X\right\}=\frac{\varrho(X, T)-\lambda}{\varrho-\lambda}
\end{gathered}
$$

where $\varrho(X, T)$ is the entropy solution of the hydrodynamic equation with initial data
$\varrho$ on the left
$\lambda$ on the right.
SSEP: CLT (of course...). Other models: interesting!

## One more result

Theorem
If $\mu^{o}$ are the stationary product marginals then, under our initial distribution, $\eta_{Q(t)}(t)$ is stationary in time.

Proof.
Repeat the argument with $\mathbf{E} \Phi\left(\eta_{i}(t)\right)$ instead of $\mathbf{E} g_{i}(t)$.

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If $\mu^{\varrho}$ are the stationary product marginals then, under our initial distribution, $\eta_{Q(t)}(t)$ is stationary in time.

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This was not even a question with exclusion.

## One more result

Theorem
If $\mu^{e}$ are the stationary product marginals then, under our initial distribution, $\eta_{Q(t)}(t)$ is stationary in time.

Proof.
Repeat the argument with $\mathbf{E} \Phi\left(\eta_{i}(t)\right)$ instead of $\mathbf{E} g_{i}(t)$.
This was not even a question with exclusion.
Only the site $Q(t)$ !

## TASEP and the corner growth model

## TASEP: Interacting particles



$$
\text { Bernoulli( } \varrho) \text { distribution }
$$

## TASEP: Interacting particles



## Bernoulli( $\varrho$ ) distribution

(particle, hole) pairs become (hole, particle) pairs with rate 1.

## TASEP: Interacting particles



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That is: waiting times $\otimes \sim$ Exponential(1).

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## TASEP: Interacting particles



## Bernoulli( $\varrho$ ) distribution

(particle, hole) pairs become (hole, particle) pairs with rate 1.
That is: waiting times $\otimes \sim$ Exponential(1).

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Particles try to jump to the right, but block each other.

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Bernoulli( $\varrho$ ) distribution
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Particles try to jump to the right, but block each other.

The Bernoulli( $\varrho$ ) distribution is time-stationary for any $(0 \leq \varrho \leq 1)$. Any translation-invariant stationary distribution is a mixture of Bernoullis.

## TASEP: Surface growth



Bernoulli( $\varrho$ ) distribution

## TASEP: Surface growth



$$
\text { Bernoulli( } \varrho) \text { distribution }
$$

## TASEP: Surface growth



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Bernoulli( $\varrho)$ distribution

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Bernoulli( $\varrho$ ) distribution

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Bernoulli( $\varrho$ ) distribution

TASEP: Last passage percolation


## TASEP: Last passage percolation



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Occupation of $(i, j)=$ jump of $P_{j}$ over $H_{i}$. Occupation of $(2,1)=$ jump of $P_{1}$ over $H_{2}$.


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The time when this happens $=: G_{i j}$.



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$P_{0}$ jumps according to a Poisson $(1-\varrho)$ process, governed by the right orange part


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Therefore:

$$
\left.\begin{array}{rl}
Q & \sim \text { Exponential }(1-\varrho) \\
& \sim \text { Exponential }(\varrho) \\
& \sim \text { Exponential }(1)
\end{array}\right\} \text { independently }
$$



The last passage model


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© starts ticking when its west neighbor becomes occupied

## The last passage model



Q Starts ticking when its west neighbor becomes occupied
©starts ticking when its south neighbor becomes occupied

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Q Starts ticking when both its west and south neighbors become occupied

## The last passage model


M. Prähofer and H. Spohn 2002

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## The competition interface



Ferrari, Martin, Pimentel (2005)
Which squares are infected via $(1,0)$ and via $(0,1)$ ?

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If it passes left of $(m, n)$, then $G_{m n}$ is not sensitive to decreasing the weights on the $j$-axis. If it passes below $(m, n)$, then $G_{m n}$ is not sensitive to decreasing the $\otimes$ weights on the $i$-axis.

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## Thank you.

