

# *Bootstrap percolation and Kinetically constrained models: time and length scales*

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# *Liquid/glass transition*

*"The deepest and most interesting unsolved problem in solid state theory is probably the theory of the nature of glass and the glass transition."* [Nobel prize P.W. Anderson]

Glasses display properties of both liquids and solids



# *Liquid/glass transition*

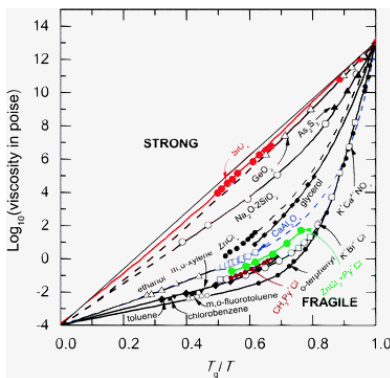
## How can you manufacture a glass?

- Take a liquid and cool it rapidly in order to prevent nucleation of the ordered crystal structure;
- relaxation times increase dramatically, the liquid falls out of equilibrium and enters a metastable phase;
- the molecules move slower and slower:  
your liquid is now a thick syrup..
- finally the liquid freezes in a structureless solid:  
here is your glass.

## *Key features of liquid/glass transition*

- huge divergence of timescales;
- no significant structural changes;
- cooperative relaxation;
- dynamical heterogeneities: non trivial spatio-temporal fluctuations, coexistence of frozen and mobile regions;
- rich phenomenology: anomalous transport properties, aging, rejuvenation, ...
- a similar jamming transition: grains in powders, emulsions, foams, colloidal suspensions, ...

# Huge relaxation times



Strong supercooled liquids: Arrhenius  $\tau \sim \exp(\Delta E/T)$

Fragile supercooled liquids: superArrhenius  $\tau \sim \exp(c/T^2), \dots$

# Kinetically Constrained Spin Models, a.k.a. KCSM

## Friedrickson Andersen model on $\mathbb{Z}^2$

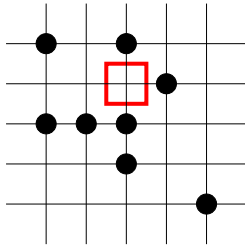
Configurations :  $\eta = \{\eta_i\}_{i \in \mathbb{Z}^2}$  with  $\eta_i \in \{0, 1\}$

Glauber dynamics = Birth and death of particles on  $\mathbb{Z}^2$

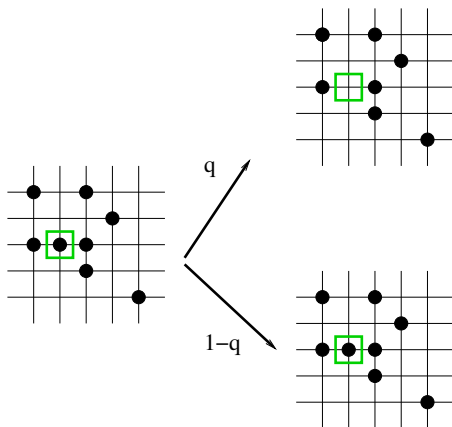
Kinetic constraint = at least 2 empty nearest neighbours

If constraint satisfied:  $1 \rightarrow 0$  rate  $q$ ,  $0 \rightarrow 1$  rate  $1 - q$

# *The kinetic constraint*



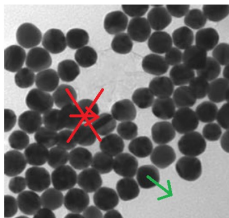
# *The kinetic constraint*





## *Ideas behind KCSM*

- Free volume shrinks when temperature is lowered;
- molecules should escape the "cage" formed by neighbours;
- local constraints act cooperatively;
- blocked structures may percolate → time scales diverge



## Other examples

- Friedrichson Andersen  $k$ -facilitated models (FA- $k$ f) on  $Z^d$ :  
at least  $k$  empty nearest neighbours
- East : at least one empty site among  $\{x - \vec{e}_1, \dots, x - \vec{e}_d\}$

# The general framework

- Choose your favorite lattice;
- Choose a collection of finite neighborhoods of 0:  
 $\{C_1, \dots, C_m\}$  with  $C_i \subset \mathbb{Z}^d \setminus 0$
- Constraint at  $x$ : at least one of the  $C_i + x$  completely empty

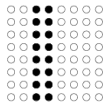
# KCSM: properties

- Constraint at  $x$  does not depend on  $\eta_x$   
→ detailed balance w.r.t. product measure

$$\mu(\eta) = \prod_{i \in \mathbb{Z}^2} q^{1-\eta_i} (1-q)^{\eta_i}$$

- $\mu$  is not the unique invariant measure
- Blocked clusters, blocked configurations

Example of blocked cluster for FA-2f and for North-East



## *Main issues, main obstacles*

- Does convergence to equilibrium at large time occurs?
- Is there a critical vacancy density below which blocked clusters percolate and relaxation is prevented?
- How does relaxation time diverge when we approach this critical density?
- **KCSM dynamics is not monotone**
- Coupling arguments and censoring not available
- Blocked configurations → relaxation not uniform on the initial condition, worst case analysis too rough
- Coercive inequalities (e.g. Log-Sobolev) anomalous

## $\mathcal{U}$ -bootstrap percolation

**Influence classes:**  $\mathcal{U} = \{C_1, \dots, C_m\}$ ,  $C_i \subset \mathbb{Z}^d$ ,  $0 \notin \cup_{i=1}^m C_i$ .

**Initial configuration**  $\eta \in \{0, 1\}^{\mathbb{Z}^d}$ .

A **deterministic discrete time** process:

- empty sites remain empty forever;
- site  $v$  is emptied at time  $t$  if the translated at  $v$  of (at least) one the influence classes  $C_i$  is completely empty at  $t - 1$ , i.e. if the **same constraint** as for **KCSM** is satisfied

Equivalent formulation:

- $A_t$  set of empty sites at time  $t$
- $A_0 := \{x \in \mathbb{Z}^d : \eta(x) = 0\}$
- $A_{t+1} := A_t \cup \{v \in \mathbb{Z}^d : v + C \subset A_t \text{ for some } C \in \mathcal{U}\}$

**Dynamics is monotone**

## Critical probability

Fix  $q \in (0, 1)$  and pick  $\eta$  random with law  $\mu = \text{Bernoulli}$  distribution with  $\mu(\eta_x = 0) = q$ .

*Does the final set of empty sites cover the lattice?*

*Which are the finite size effects ?*

Consider the process on the torus  $\mathbb{Z}_n^d$ .

$$q_c(n, \mathcal{U}) := \inf\{q \in [0, 1] : \mu(\cup_{t \geq 0} A_t = \mathbb{Z}_n^d) \geq 1/2\}$$

*How does  $q_c(n, \mathcal{U})$  depend on  $\mathcal{U}$ ? How does it scale for  $n \rightarrow \infty$ ?*

# KCSM and $\mathcal{U}$ -bootstrap

Final set of 1's for bootstrap  $\leftrightarrow$  blocked particles for KCSM

- $\mu$  is mixing for a KCSM iff  $q > q_c := \liminf_{n \rightarrow \infty} q_c(n, \mathcal{U})$
- How fast do we converge to  $\mu$ ? Exponentially  
 $\forall q > q_c, \exists T_{rel}(q) < \infty$  s.t.  
 $\mu(fP_t g) - \mu(f)\mu(g) \leq C_{f,g} \exp(-t/T_{rel}(q)), \quad \forall f, g \in L^2(\mu)$
- How does  $T_{rel}$  (=inverse of spectral gap) diverge as  $q \downarrow q_c$ ?  
Set  $L_c(q) := \min\{n : q_c(n, \mathcal{U}) = q\}$ . Then

$$L_c \leq T_{rel} \leq e^{L_c^d}$$

[Cancrini, Martinelli, Roberto, C.T. '08]



## Relaxation time for FA-kf model, $k \leq d$

$q_c = \liminf_{n \rightarrow \infty} q_c(n, \mathcal{U}) = 0$  [Van Enter '87, Schonmann '90]

$$\exists \lambda(d, k) > 0 \text{ s.t. } L_c = \exp_{k-1} \left( \frac{\lambda(d, k) + o(1)}{q^{1/(d-k+1)}} \right)$$

[Aizenmann, Lebowitz '88, Cerf, Manzo '02, Balogh, ..., Bollobas, Duminil-Copin, Morris '12]

### Theorem (Martinelli, C.T. '16)

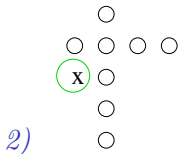
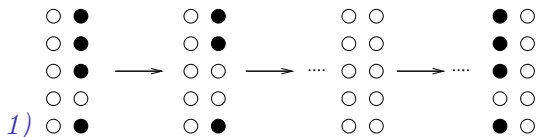
- For the FA-2f model there exists  $\alpha, c > 0$  s.t.

$$\exp(c/q^{1/(d-1)}) \leq T_{rel} \leq \exp \left( \log(1/q)^\alpha / q^{1/(d-1)} \right)$$

- For the FA-kf model for any  $k \geq 3$  there exists  $c, c'$  s.t.

$$\exp_{k-1} \left( \frac{c}{q^{1/(d-k+1)}} \right) \leq T_{rel} \leq \exp_{k-1} \left( \frac{c'}{q^{1/(d-k+1)}} \right)$$

# FA-2f Dominant relaxation mechanism



3) Set  $\ell_q := 1/q \log 1/q$ .

$\mu(\text{a segment of length } \ell_q \text{ contains at least one vacancy}) \sim 1$

# $FA - 2f$ $d = 2$ , strategy of the proof

$$T_{rel} := \inf\{\lambda : Var(f) \leq \lambda \mu(f, -\mathcal{L}f)\} \quad \forall f \text{ local}$$

$$\mu(f, -\mathcal{L}f) = \sum_{x \in \mathbb{Z}^2} \mu(c_x Var_x(f))$$

$c_x$  = indicator function that  $x$  has  $\geq 2$  empty nearest neighb.

$Var_x$  = local variance at  $x$

We want to prove  $T_{rel} \leq \exp(c|\log q|/q)$ , i.e. that  $\forall f$  it holds

$$Var(f) \leq \exp(c|\log q|/q) \sum_{x \in \mathbb{Z}^2} \mu(c_x Var_x(f))$$

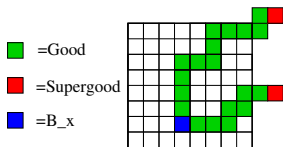
## Step 1: a key constrained Poincaré inequality

Renormalize on  $L \times L$  boxes with  $L = 1/q \log 1/q$ .

- a box is **good** if it contains at least one empty site on each column and on each line  $\rightarrow \mu(\text{good}) \sim 1$
- a box is **super good** if it is good and contains at least one empty column and one empty row  $\rightarrow \mu(\text{super good}) \sim \exp(-1/q \log(1/q)^2) \ll 1$

$$\text{Var}(f) \leq \sum_{x \in \mathbb{Z}^2(L)} \mu(\kappa_x \text{Var}_{B_x})$$

$\kappa_x$  = indicator function of

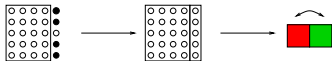


Maximal path length  
 $\exp(1/q \log(1/q)^2)$

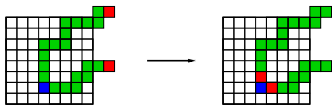
Very flexible tool

## Step 2: construct allowed paths

Swap neighbouring good and supergood boxes



Bring two supergood boxes near  $B_x$



For any  $\omega \in \{0, 1\}^{B_x}$  and  $y \in B_x$  we can now bring an empty row and column near  $y \in B_x$ : **the constraint at  $y$  is now satisfied**



## Step 3: canonical paths for reversible Markov chains

Key ingredient: the whole path is constructed by "shifting" an empty column of height  $L = 1/q \log 1/q$

Our constrained Poincaré inequality + canonical paths for reversible Markov chains

$$\rightarrow T_{rel} \leq \exp(c/q(\log(1/q))^2)$$

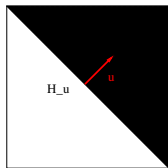
(= length of the path  $\times$  congestion constant)

Changing the notion of Good, Supergood, L, and the oriented neighborhood of  $B_x$  we cover other models..all critical models?

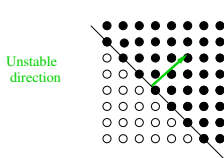
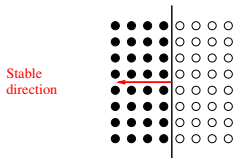
# A universality result for cellular automata in $\mathbb{Z}^2$

Take  $u \in S^1$ , let  $H_u := \{x \in \mathbb{Z}^2 : \langle x, u \rangle < 0\}$ .

$u$  is a **stable direction** if starting from  $\eta$  empty on  $H_u$  and filled on  $\mathbb{Z}^2 \setminus H_u$  no other site can be emptied.



Ex. East:  $\vec{u} = -\vec{e}_1$  is **stable**;  $\vec{u} = \vec{e}_1 + \vec{e}_2$  is **unstable**



## Classification of cellular automata in $\mathbb{Z}^2$

- **supercritical** if  $\exists$  open semicircle without stable directions;
- **critical** if every open semicircle has a stable direction and  $\exists$  a semicircle with a finite number of stable directions
- **subcritical** otherwise

Red= stable direction: Green= unstable direction

FA1f : supercritical



S

East: supercritical



S

FA2f: critical



S

North-East : subcritical



S



# A universality result for cellular automata in $\mathbb{Z}^2$

*Theorem [Bollobas, Smith, Uzzell '15 + Bollobas, Duminił-Copin, Morris, Smith '16 + Balister, Bollobas, Przykucki, Smith '16]*

- Supercritical models:  $q_c(n, \mathcal{U}) = (1/n)^{\Theta(1)}$
- Critical models:  $\exists \alpha(\mathcal{U}) > 0$  s.t.  $q_c(n, \mathcal{U}) = \Theta(1/\log n)^\alpha$
- Subcritical models:  $\liminf_{n \rightarrow \infty} q_c(n, \mathcal{U}) > 0$

$\rightarrow L_c(q) = 1/q^{\Theta(1)}$  for supercritical models

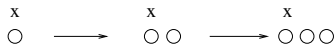
$\rightarrow L_c = \Theta(\exp(1/q^\alpha))$  for critical models

*$L_c(q)$  determined by the action of the cellular automata on discrete half planes*

# Supercritical models the key mechanism

## Supercritical models:

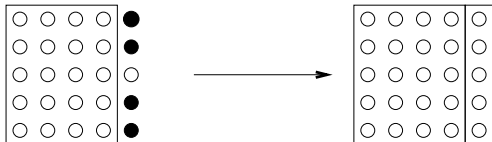
- there is a finite empty droplet  $D \subset \mathbb{Z}^d$  from which we can empty an infinite half line
- if  $n \gg 1/q^{|D|}$  we will typically droplets to empty all  $\mathbb{Z}_n^d$
- Ex. East: a single empty site is a droplet



# Critical models the key mechanism

## Critical models:

- we can expand a finite empty droplet one step further iff we find a group of  $\alpha$  empty sites on its boundary
- if we have an empty droplet of size  $\gg 1/q^\alpha$  we will typically be able to continue until emptying all  $\mathbb{Z}_n^d$ .
- Ex. FA-2f: a rectangle of empty sites can be expanded if there is at least one empty on the next column



# Supercritical KCSM on $\mathbb{Z}^2$

*Theorema [Martinelli, Morris, C.T. '16]*

A refined classification : a supercritical model is **rooted** if there are two non opposite stable directions. It is **unrooted** otherwise.

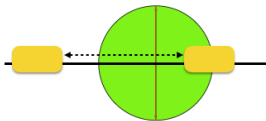
- for all supercritical unrooted models  $T_{rel} = 1/q^{\Theta(1)}$
- for all supercritical rooted models  $T_{rel} = 1/q^{\Theta(\log(1/q))}$

$$L_c = \frac{1}{q}^{\Theta(1)}$$

→ unrooted models  $\exists \alpha$  s.t.  $T_{rel} = O(L_c^\alpha)$

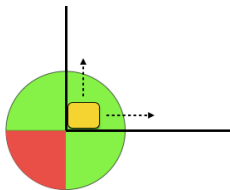
→ rooted models  $\exists \alpha$  s.t.  $T_{rel} \geq L_c^{\alpha \log L_c}$

## Intuition behind the unrooted result



- $\exists$  empty droplet can be shifted back and forth along a line: from the droplet one can empty the entire line
- FA-1 f is unrooted, empty droplet = a single empty site
- scaling proven via renormalization to FA1f model in  $d = 1$  and using the polynomial result for FA1f [Cancrini, Martinelli, Roberto, Toninelli '08]

## Intuition behind the rooted result



- from any finite empty region we can empty only a cone.
- lower bound: **logarithmic energy barriers**
  - start from a single droplet
  - to create a new droplet at distance  $\ell$  you necessarily go through a configuration with  $c \log \ell$  simultaneous empty sites
  - $\rightarrow$  time  $1/q^{c \log 1/q}$
- upper bound: **renormalisation to East** in  $d = 1$  model and using  $T_{rel}^{East} = 1/q^{c \log 1/q}$  [Aldous, Diaconis '02, Cancrini, Martinelli, Roberto, Toninelli '08, Chlebloun, Faggionato, Martinelli '15]

## Critical KCSM on $\mathbb{Z}^2$

$\alpha(\vec{u}) =$  **difficulty of direction  $\vec{u}$**  = minimal number of empty site to be added to  $H_u$  in order to grow the empty set  $H_u$  of one step in the  $\vec{u}$  direction

$$\alpha := \min_C \max_{\vec{u} \in C} \alpha(\vec{u})$$

$$\rightarrow L_c = \Theta(\exp(c/q^\alpha))$$

### Conjecture

[Martinelli, Morris ,C.T.] A refined classification: a critical model is  **$\alpha$ -rooted** if there are two non opposite stable directions of difficulty  $> \alpha$ . It is  **$\alpha$ -unrooted** otherwise. Then

- for  $\alpha$ - unrooted models  $T_{rel} = O(\exp(c/q^\alpha |\log(1/q)|))$
- for  $\alpha$ -rooted models  $T_{rel} \geq \exp(c/q^\beta)$ ,  $\beta > \alpha$

## Summary

- KCSM are stochastic models for liquid/glass transition
- intimate relation to bootstrap percolation
- ergodicity for KCSM = percolation cellular automata
- $T_{rel} = 1/\text{gap} < \infty$  in the ergodic regime and  $T_{rel} > L_c$
- universality results for bootstrap percolation in  $d = 2$
- due to logarithmic barriers sometimes  $T_{rel} \gg L_c$
- a refined classification of the influence classes, conjecture on the universal behavior for supercritical / critical KCSM
- a new toolbox to upper bound  $T_{rel}$ 
  - scaling for FA- $kf$  on  $\mathbb{Z}^d$
  - scaling for all critical/supercritical models in  $d = 2$  hopefully ...