Connecting blocking measures of interacting particle systems with combinatorial objects

(Based on joint works with Daniel Adams, Márton Balázs, Dan Fretwell and Benjamin Lees)

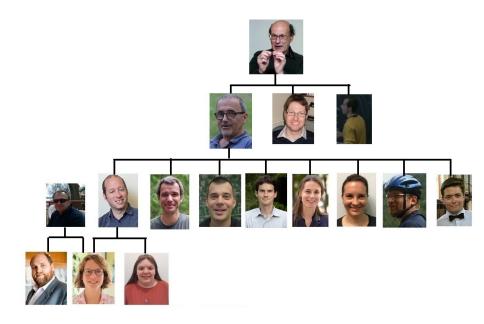
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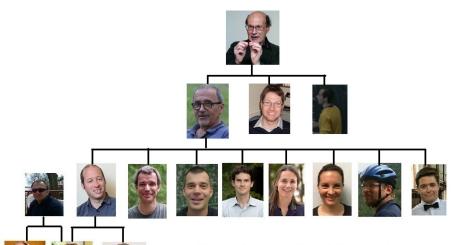
Stochastics and Influences Workshop On the ocassion of Bálint Tóth's 70th Birthday



Connecting blocking measures of interacting particle systems with combinatorial objects

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I have a truly marvellous full academic tree which this slide is too small to contain!

Probability and Combinatorics

We all know that Probability and Combinatorics are closely related! Connecting blocking measures of interacting particle systems with combinatorial objects

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Blocking Ising

Probability and Combinatorics

We all know that Probability and Combinatorics are closely related!

For example integrable systems and young tableaux

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Probability and Combinatorics

We all know that Probability and Combinatorics are closely related!

For example integrable systems and young tableaux

 Recently Blocking Measures for Interacting Particle Systems have been shown to be linked to combinatorial objects. Connecting blocking measures of interacting particle systems with combinatorial objects

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Balázs and Bowen (2018) gave a purely probabilistic proof of a well-known combinatorial identity: Connecting blocking measures of interacting particle systems with combinatorial objects

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Balázs and Bowen (2018) gave a purely probabilistic proof of a well-known combinatorial identity:

Jacobi Triple Product

For $q \in (0,1)$ and $z \neq 0$,

$$\sum_{m\in\mathbb{Z}}q^{rac{m(m+1)}{2}}z^m=\prod_{i\geq 1}(1-q^i)(1+q^iz)(1+q^{i-1}z^{-1})$$

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This paper constructs a family of blocking measures for Interacting Particle Systems on Z. Connecting blocking measures of interacting particle systems with combinatorial objects

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- This paper constructs a family of blocking measures for Interacting Particle Systems on Z.
- This proof of the Jacobi Triple Product identity follows from the Exclusion - Zero-range correspondence

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The story begins ...

Balázs and Bowen (2018) gave a purely probabilistic proof of a well-known combinatorial identity:

Jacobi Triple Product (JTP)

For $q \in (0,1)$ and $c \in \mathbb{R}$,

$$rac{\sum\limits_{m \in \mathbb{Z}} q^{rac{m(m+1)}{2} - mc}}{\prod\limits_{i \geq 1} (1 + q^{i-c})(1 + q^{i-1+c})} = \prod\limits_{i \geq 1} (1 - q^i)$$

- This paper constructs a family of blocking measures for Interacting Particle Systems on Z.
- This proof of the Jacobi Triple Product identity follows from the Exclusion - Zero-range correspondence, by equating blocking measures.

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The story continues ...

Natural questions for ASEP leads to proofs of other well known combinatorial identities. (Adams, Balázs, J.)

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Distribution of #{particles to left of a site} leads to,

Euler's Identity For $q \in (0,1)$ and $z \in \mathbb{R}_{>0}$, $\sum_{k=0}^{\infty} \frac{q^{\frac{k(k-1)}{2}z^k}}{\prod\limits_{i=1}^k (1-q^i)} = \prod_{i=0}^{\infty} (1+zq^i).$

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Euler's Identity
For
$$q \in (0,1)$$
 and $z \in \mathbb{R}_{>0}$, $\sum_{k=0}^{\infty} \frac{q^{\frac{2}{2}}z^k}{\prod\limits_{i=1}^k (1-q^i)} = \prod_{i=0}^{\infty} (1+zq^i).$

Particle - Hole symmetry for ASEP leads to,

Durfee Rectangles Identity

For $q \in (0,1)$ and any fixed $n \in \mathbb{Z}$,

$$\frac{1}{\prod\limits_{i\geq 1} (1-q^i)} = \sum_{k=\max\{-n,0\}}^{\infty} \frac{q^{k(n+k)}}{\prod\limits_{i=1}^{n+k} (1-q^i) \cdot \prod\limits_{j=1}^{k} (1-q^j)}.$$

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Natural Questions:

Are ASEP/AZRP the only particle systems that have this connection to combinatorics via these identities?

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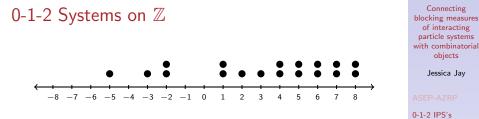
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Natural Questions:

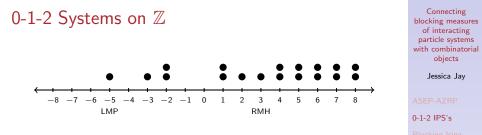
Are ASEP/AZRP the only particle systems that have this connection to combinatorics via these identities?

If this connection is deeper, does it reveal new results in Probability or Combinatorics? Connecting blocking measures of interacting particle systems with combinatorial objects

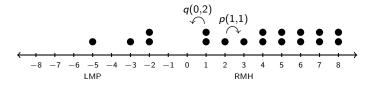
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0-1-2 Systems on $\mathbb Z$



Jump rate functions: Right jump from site *i* to i + 1 w/r $p(\eta_i, \eta_{i+1})$. Left jump from site *i* to i - 1 w/r $q(\eta_{i-1}, \eta_i)$.

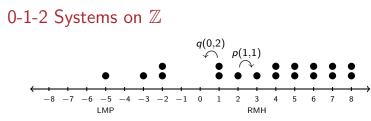
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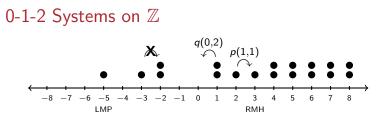


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• Exclusion: $p(0, \cdot) = p(\cdot, 2) = q(2, \cdot) = q(\cdot, 0) = 0$

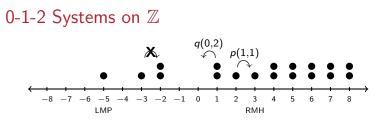
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Blocking Ising



- Exclusion: $p(0, \cdot) = p(\cdot, 2) = q(2, \cdot) = q(\cdot, 0) = 0$
- Attractivity: $p(y+1,z) \ge p(y,z), \ p(y,z+1) \le p(y,z), \ q(y+1,z) \le q(y,z), \ q(y,z+1) \ge q(y,z)$

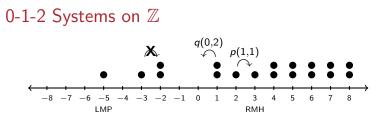
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• Attractivity:
$$p(y+1,z) \ge p(y,z)$$
, $p(y,z+1) \le p(y,z)$,
 $q(y+1,z) \le q(y,z)$, $q(y,z+1) \ge q(y,z)$

• Algebraic recursive rates: an explicit condition. The rates can be parametrised by $q \in (0, 1)$ and $t \ge 1$,

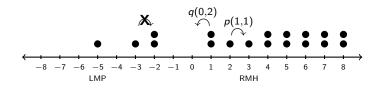
$$q:=rac{q(0,1)}{p(1,0)}, \qquad t:=\left(rac{p(1,0)q(0,2)}{q(0,1)p(1,1)}
ight)^{rac{1}{2}}.$$

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0-1-2 Systems on $\ensuremath{\mathbb{Z}}$



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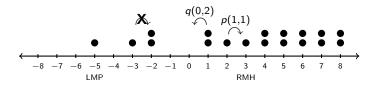
2-Lane ASEP

Family of reversible stationary measures

For any $c \in \mathbb{R}$ the following is reversible and stationary for the dynamics,

$$\underline{\mu}^{c}(\underline{\eta}) = \prod_{i=-\infty}^{0} \frac{t^{\mathbb{I}\{\eta_{i}=1\}} q^{(c-i)\eta_{i}}}{(1+tq^{c-i}+q^{2(c-i)})} \prod_{i=1}^{\infty} \frac{t^{\mathbb{I}\{\eta_{i}=1\}} q^{(2-\eta_{i})(i-c)}}{(1+tq^{i-c}+q^{2(i-c)})}$$

0-1-2 Systems on $\ensuremath{\mathbb{Z}}$



Conserved Quantity: $N(\underline{\eta}) = \sum_{i=1}^{\infty} (2 - \eta_i) - \sum_{i=-\infty}^{0} \eta_i < \infty$

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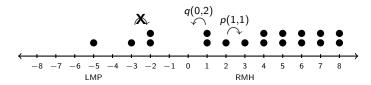
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0-1-2 Systems on $\ensuremath{\mathbb{Z}}$



Conserved Quantity:
$$N(\underline{\eta}) = \sum_{i=1}^{\infty} (2 - \eta_i) - \sum_{i=-\infty}^{0} \eta_i < \infty$$

Unique stationary measures

The unique stationary measure on $\{\underline{\eta}: N(\underline{\eta}) = n\}$ (for *n* even) is given by,

$$\underline{\nu^{n}}(\underline{\eta}) = \frac{2\sum_{\ell=-\infty}^{\infty} q^{\ell(\ell+1)-2\ell c} \underline{\mu}^{c}(\underline{\eta}) \mathbb{I}\{N(\underline{\eta}) = n\}}{q^{\frac{n(n+2)}{4}-nc} \left(1 + \prod_{i=-\infty}^{\infty} \left(1 - 2\mu_{i}^{c}(1)\right)\right)}$$

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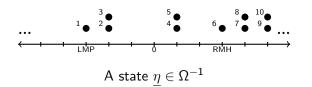
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For each $n \in \mathbb{Z}$ there is a bijection, $T^n : \Omega^n \to \mathcal{H}$, for some specific state space $\mathcal{H} \subset \mathbb{Z}_{\geq 0}^{\mathbb{Z}_{< 0}}$.



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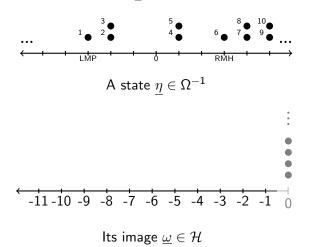
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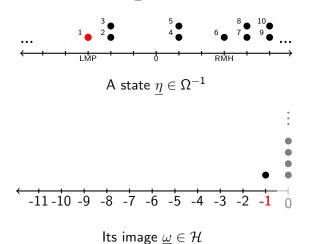
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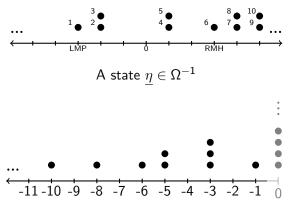
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Its image $\underline{\omega} \in \mathcal{H}$

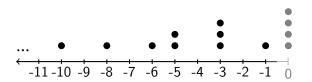
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State space:

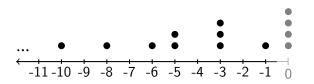
• Any number of particles allowed per site, $\mathcal{H} \subset \mathbb{Z}_{>0}^{\mathbb{Z}_{<0}}$

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State space:

- Any number of particles allowed per site, $\mathcal{H} \subset \mathbb{Z}_{>0}^{\mathbb{Z}_{<0}}$
- Restriction: No 2 consecutive empty sites

 $\mathcal{H} \subset \mathcal{H}' := \{ \underline{\omega} \in \mathbb{Z}_{\geq 0}^{\mathbb{Z}_{< 0}}: \ \omega_{-i} = 0 \Rightarrow \omega_{-i-1} \neq 0, \forall i > 0 \}$

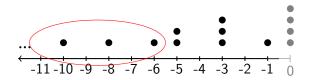
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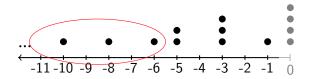
► Far to the left agrees with 1,0,1,0,1... $\mathcal{H} = \mathcal{H}^e \cup \mathcal{H}^o$ $\mathcal{H}^e := \{ \underline{\omega} \in \mathcal{H}' : \exists N > 0 \text{ s.t } \omega_{-i} = \mathbb{I}\{i \text{ even}\} \ \forall i \geq N\},\$ $\mathcal{H}^o := \{ \underline{\omega} \in \mathcal{H}' : \exists N > 0 \text{ s.t } \omega_{-i} = \mathbb{I}\{i \text{ odd}\} \ \forall i \geq N\}.$ Connecting blocking measures of interacting particle systems with combinatorial objects

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Blocking Ising

To find the reversible stationary measures on \mathcal{H}^e and \mathcal{H}^o ,

First relax the restriction on the dynamics and use Balázs and Bowen blocking measure construction. Connecting blocking measures of interacting particle systems with combinatorial objects

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Blocking Ising

To find the reversible stationary measures on \mathcal{H}^e and \mathcal{H}^o ,

- First relax the restriction on the dynamics and use Balázs and Bowen blocking measure construction.
- Markov "cuts" Theorem gives that this is also reversible stationary for the process with restricted dynamics.

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Blocking Ising

Family of Restricted Particle Systems

To find the reversible stationary measures on \mathcal{H}^e and $\mathcal{H}^o,$

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- Condition on having the 1,0,1,0,1,0 ... block far to the left (this has measure zero but can use the product structure of the measure)

Reversible stationary measure

On \mathcal{H}^e the **unique reversible stationary** measure is,

$$\underline{\pi}^{e}(\underline{\omega}) = \frac{q^{\sum\limits_{i \text{ odd}} i\omega_{-i} + \sum\limits_{i \text{ even}} i(\omega_{-i}-1)} t^{2\left(\sum\limits_{i \text{ odd}} \mathbb{I}\{\omega_{-i} \ge 1\} - \sum\limits_{i \text{ even}} \mathbb{I}\{\omega_{-i} = 0\}\right)}{S_{\text{even}}(q, t)}$$

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This equivalence proves new Jacobi style identities!

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This equivalence proves new Jacobi style identities!

Theorem (Balázs, Fretwell, J., (2022))
For
$$0 < q < 1$$
, $t \ge 1$ and $z > 0$
 $2S_{even}(q, t) \sum_{\ell \in \mathbb{Z}} q^{\ell(\ell+1)} z^{2\ell}$
 $= \prod_{i\ge 1} (1 + tzq^i + z^2q^{2i})(1 + tz^{-1}q^{i-1} + z^{-2}q^{2(i-1)})$
 $+ \prod_{i\ge 1} (1 - tzq^i + z^2q^{2i})(1 - tz^{-1}q^{i-1} + z^{-2}q^{2(i-1)})$

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Theorem (Balázs, Fretwell, J., (2022)) For 0 < q < 1, t > 1 and z > 0 $2S_{\rm even}(q,t)\sum q^{\ell(\ell+1)}z^{2\ell}$ $= \prod (1 + tzq^{i} + z^{2}q^{2i})(1 + tz^{-1}q^{i-1} + z^{-2}q^{2(i-1)})$ $i \geq 1$ + $\prod (1 - tzq^{i} + z^{2}q^{2i})(1 - tz^{-1}q^{i-1} + z^{-2}q^{2(i-1)})$ $i \geq 1$

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This identity was found from the probability and then we found the combinatorial interpretation.

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When we specialise to the 2-Exclusion process (all possible left jumps with rate q, right with rate 1) the identity specialises to a known one in the combinatorics literature.

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$$\mathcal{H}^{e} = \{ \underline{\omega} \in \mathbb{Z}_{\geq 0}^{\mathbb{Z}_{<0}} : \ \omega_{-i} = 0 \Rightarrow \omega_{-i-1} \neq 0, \forall i > 0$$

and $\exists N > 0 \quad \text{s.t } \omega_{-i} = \omega_{-i}^{e} \ \forall i \geq N \}$

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and $\exists N > 0 \quad \text{s.t } \omega_{-i} = \omega_{-i}^{e} \ \forall i \geq N \}$

Combinatorics gives:

$$\begin{split} &\sum_{\underline{\omega}\in\mathcal{H}^{e}}q^{\sum\limits_{i\text{ odd}}i\omega_{-i}+\sum\limits_{i\text{ even}}i(\omega_{-i}-1)}\\ &=\frac{1}{\prod\limits_{i\geq 1}(1-q^{i})(1-q^{12i-10})(1-q^{12i-9})(1-q^{12i-3})(1-q^{12i-2})} \end{split}$$

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and $\exists N > 0 \quad \text{s.t } \omega_{-i} = \omega_{-i}^{e} \ \forall i \geq N \}$

Combinatorics gives:

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This is completely non-obvious from the form of the blocking measure!

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0-1-2 IPS's

Blocking Ising

A mystery remains ...

Connecting blocking measures of interacting particle systems with combinatorial objects

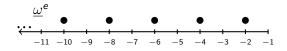
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$$\mathcal{H}^{e} = \{ \underline{\omega} \in \mathbb{Z}_{\geq 0}^{\mathbb{Z}_{< 0}} : \omega_{-i} = 0 \Rightarrow \omega_{-i-1} \neq 0, \forall i > 0$$

and $\exists N > 0 \quad \text{s.t } \omega_{-i} = \omega_{-i}^{e} \ \forall i \geq N \}$

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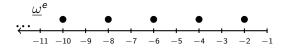
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Conjecture:

$$\sum_{\omega \in \mathcal{H}^{e}} q^{\sum_{i \text{ odd}} i\omega_{-i} + \sum_{i \text{ even}} i(\omega_{-i}-1)} t^{2\left(\sum_{i \text{ odd}} \mathbb{I}\{\omega_{-i} \ge 1\} - \sum_{i \text{ even}} \mathbb{I}\{\omega_{-i}=0\}\right)} \\ = \frac{1}{\prod_{m \ge 1} (1-q^{2m})} + \frac{\sum_{i \ge 1} \sum_{n \ge i} (-1)^{n-i} \binom{n+i-1}{2i-1} \frac{n}{i} q^{n^{2}} t^{2i}}{\prod_{m \ge 1} (1-q^{m})^{2}}$$

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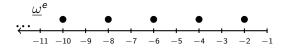
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This is not clear from either the probability or combinatorics interpretation of our identities.

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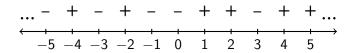
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Blocking Ising

(Blocking) Ising process State space: $\Omega^{Is} := \{-1, +1\}^{\mathbb{Z}}$



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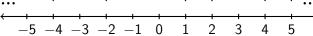
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To any state we associate two quantities:

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(Blocking) Ising process State space: $\Omega^{ls} := \{-1, +1\}^{\mathbb{Z}}$... - + - + - - + + - + + + +

To any state we associate two quantities:

• "Energy":
$$H(\sigma) := \sum_{i \in \mathbb{Z}} \mathbb{I}\{\sigma_i \neq \sigma_{i+1}\}.$$

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Blocking Ising

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$$H(\sigma) := \sum_{i \in \mathbb{Z}} \mathbb{I}\{\sigma_i \neq \sigma_{i+1}\}.$$

• "Asymmetry": for any given $c \in \mathbb{R}$,

$$f_c(\sigma) := 2 \sum_{i=1}^{\infty} (i-c) \mathbb{I}\{\sigma_i = -1\} - 2 \sum_{i=-\infty}^{0} (i-c) \mathbb{I}\{\sigma_i = 1\}.$$

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We can also consider inhomogeneous interactions

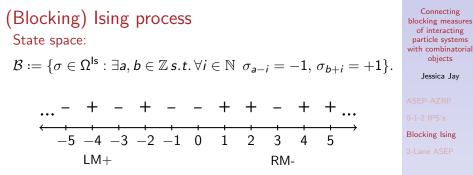
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To any state we associate two quantities:

"Energy":

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We consider Kawasaki dynamics for the Ising process, where spins swap at sites (i, i + 1) with rates of the form,

$$w(\sigma,\sigma')=rac{1}{2}(1\pm anh(etarac{J(i-1)\pm J(i+1)}{2}))q^{\pm 1}.$$

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Compared to ASEP: the rate of a jump over edge (i, i + 1) now also depends on sites i - 1 and i + 2.

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Family of Stationary Reversible Measures

For any $c \in \mathbb{R}$ the measure,

$$\mu^{c}(\sigma) := \frac{e^{-\beta H(\sigma)}q^{f_{c}(\sigma)}}{Z_{\beta,q,c}}$$

is stationary and reversible for the described dynamics.

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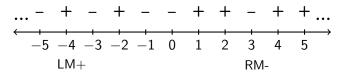
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Conserved Quantity:

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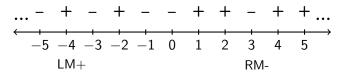
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Unique stationary measures

The unique stationary measure on $\{\sigma : N(\sigma) = n\}$ is,

$$\nu_{\beta,q}^{n}(\sigma) = \frac{\sum\limits_{m \in \mathbb{Z}} q^{m(m+1)-2mc} \cdot e^{-\beta H(\sigma)} q^{f_{c}(\sigma)}}{q^{n(n+1)-2nc} Z_{\beta,q,c}} \mathbb{I}\{N(\sigma) = n\}$$

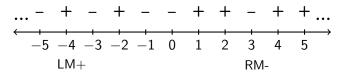
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Blocking Ising

 Using the maps from the ASEP-AZRP correspondence we can find an equivalent particle system. Connecting blocking measures of interacting particle systems with combinatorial objects

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- Using the maps from the ASEP-AZRP correspondence we can find an equivalent particle system.
- We find the unique reversible measures for this equivalent system and hence have an identity.

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- As a consequence of this identity and the Jacobi Triple Product identity we can write the partition function of the blocking Ising as a product,

$$Z_{eta,q,c} = e^{-eta} \prod_{i=1}^\infty (1\!+\!(e^{-2eta}\!-\!1)q^{2i})(1\!+\!q^{2(i-c)})(1\!+\!q^{2(i-1+c)})$$

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Without the identity/ combinatorics the partition function is very complicated!

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Blocking Ising

If we let $Q := q^2$, $z := q^{-2c}$ and $y := e^{-2\beta}$ the partition function of Blocking Ising is,

$$y + \sum_{\substack{L,R \ge 0 \\ L+R > 0}} y^{(L+R+1)} \sum_{\substack{\ell_1, \dots, \ell_L \ge 1 \\ m_1, \dots, m_R \ge 1}} \prod_{j=1}^L \frac{Q^{\frac{1}{2}\ell_j(\ell_j-1) + (\ell_{L-j}+1)(\ell_L + \dots + \ell_{L-j+1})_z - \ell_j}}{1 - Q^{\ell_L + \dots + \ell_{L-j+1}}}$$

$$\prod_{j=1}^{R} \frac{Q^{\frac{1}{2}m_{j}(m_{j}-1)+(m_{R-j}+1)(m_{R}+\cdots+m_{R-j+1})}z^{m_{j}}}{1-Q^{m_{R}+\cdots+m_{R-j+1}}} \left((1-y^{-1})Q^{\ell_{L}+\cdots+\ell_{1}}Q^{m_{R}+\cdots+m_{1}}+y^{-1} \right).$$

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So the combinatorial identity gives us something highly non trivial!

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Amir, Bahadoran, Busani, Saada (2023) characterised invariant measures for multi-lane exclusion (including blocking measures). Connecting blocking measures of interacting particle systems with combinatorial objects

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We are seeing identities related to the ones from 0-1-2 systems. Connecting blocking measures of interacting particle systems with combinatorial objects

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- We are seeing identities related to the ones from 0-1-2 systems.
- The parameter t in the 0-1-2 system can be seen as some sort of strength of gravity.
- We believe we can extend Balázs and Bowen's blocking family to 2 (potentially multiple) lanes.

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Often these combinatorial identities also have algebraic meaning!

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- Often these combinatorial identities also have algebraic meaning!
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- Question: Can these particle system equivalences be seen algebraically?
- Question: Can the link with algebra give us more new results in probability?

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Boldog Születésnapot Bálint!



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